

# Building and observing jets in and out the medium

Gregory Soyez

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CERN — Joint Heavy-Ion theory & QCD

I do not consider myself a heavy-ion physicist

- My background is more on (jets from) the “vacuum” side of high-energy collisions
- The pure “heavy-ion” part of this talk is most likely biased towards my own work with Paul Caucal and Edmond Iancu

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## I have mixed feelings about the very existence of this talk

- After all, heavy-ion collisions are QCD collisions!  
What bridges are we building then?
- Conversely, HI collisions are substantially more complex  
⇒ more “modelling” & qualitative arguments

# From LEP/RHIC/Tevatron to the LHC (quite obvious)

## 1 Higher energy

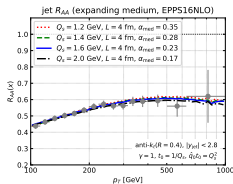
Main consequence:  
large pQCD  
phase-space

- ✓ precision physics
- ✓ th. uncertainties

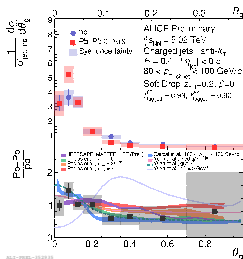
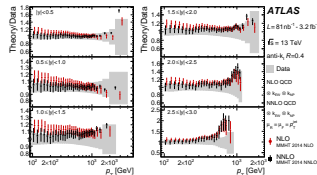
## 2 Higher luminosity

- ✓ more observables
- ✓ higher precision
- ✓ more differential

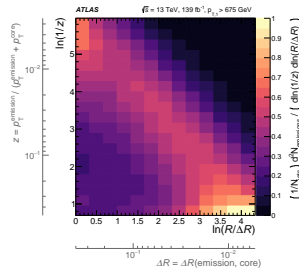
PC, EI, GS, 2012.01457



ATLAS, 1711.02692



ALI-PREL-352935



ATLAS, 2004.03540

# Different physics at different scales: $pp$

LHC probes physics across many scales

$Q \equiv$   
100 GeV  
or higher

$\mu_{\text{NP}} \sim$   
1 GeV

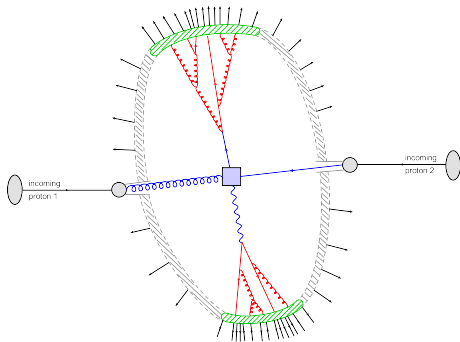
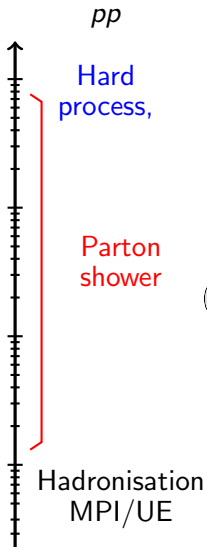


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$pp$

Hard  
process,

Parton  
shower

Hadronisation  
MPI/UE

A lot of work in past 20 years:

- amplitudes (NLO, NNLO)
- NLO availability: MadGraph, aMC@NLO, POWHEG, MCFM, ...
- matching/merging: MLM, CKKW Mi(N)NLO, UNNLOPS, Geneva, ...
- $N^{k \geq 2}$  LL resummations
- Historical showers: Pythia, Herwig, Sherpa
- More recent work: Dire, Vincia, Deductor, PanScales

Nonperturbative modelling

$$\propto (\mu_{NP}/Q)^{\#}$$

if IRC-safe

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some still need to make it to HI

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# Different physics at different scales: HI

Jets scatter off the QGP

$k_t$  broadening

Transverse “kicks”:

$$\hat{q} \equiv \langle k_{\perp}^2 \rangle \text{ per unit length.}$$

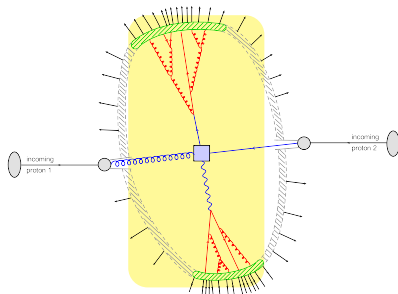
Relevant scale:  $Q_s = \sqrt{\hat{q}L}$

Medium-induced emissions

~ QCD colour coherence broken  
⇒ extra emissions (mostly large-angle)

2 relevant scales:

- $\hat{q}L^2$ : hard (rare) emissions
- $\alpha_s^2 \hat{q}L^2$ : multiple emissions



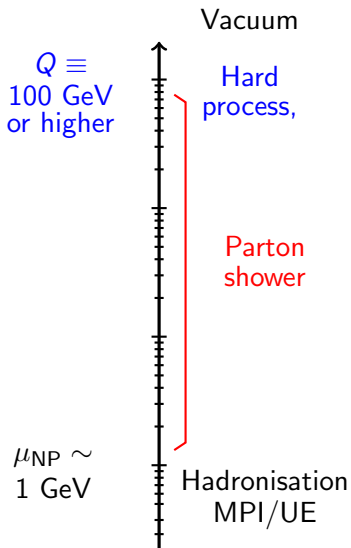
effects on the QGP itself

No real separation between the jet and the medium  
⇒ correlated behaviour  
back-reaction, medium recoil, ...

Relevant scale  $\sim T$

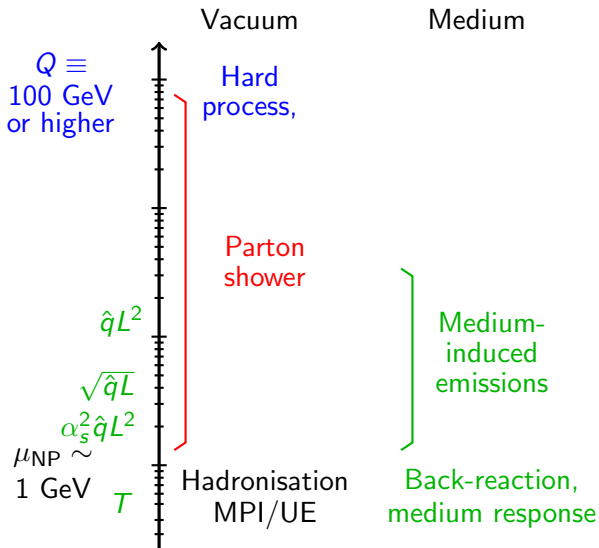
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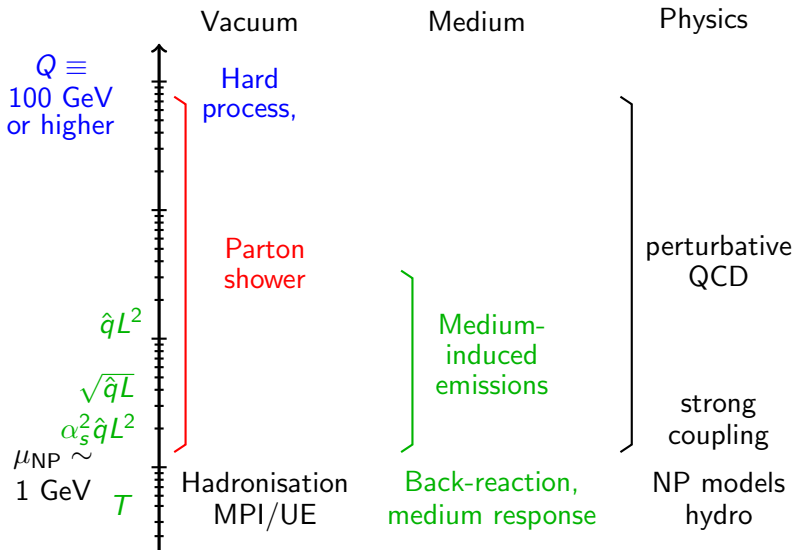
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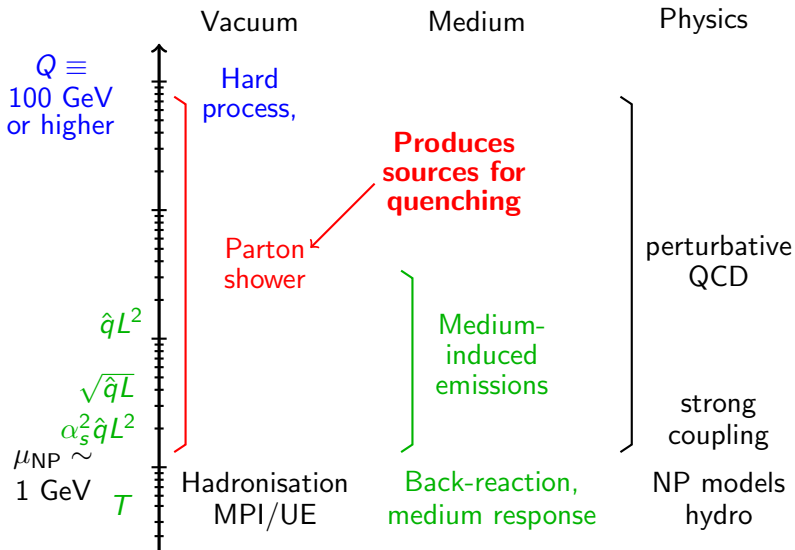
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# Substructure opens (almost) endless options

**Large phase-space where perturbative QCD can genuinely be applied both in  $pp$  and in HI**

implies correlated effort on (at least) 2 fronts

Understand the development  
of parton cascades

develop observables sensitive  
to the QCD dynamics

Calls for several improvements:

- fixed-order ( $pp$ )
- resummations ( $pp$ )
- quenching description (HI)
- parton-shower MC (both)

Jet substructure has proven a powerful tool for about 10 years

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Jet substructure has proven a powerful tool for about 10 years

I will focus on these 2 topics with deep connections between  $pp$  and HI

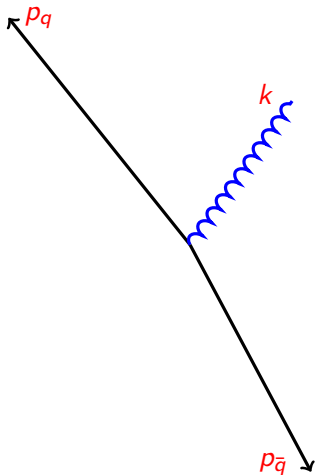
# Interlude

## Representing radiation in the Lund plane



# Radiation from a $q\bar{q}$ antenna

Take a gluon emission from a  $(q\bar{q})$  dipole



Emission:

$$k^\mu \equiv z_q p_q^\mu + z_{\bar{q}} p_{\bar{q}}^\mu + k_\perp^\mu$$

3 degrees of freedom:

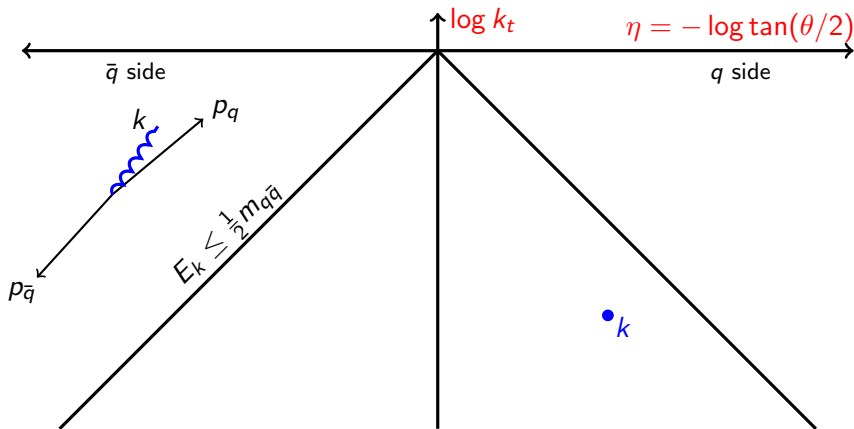
- Rapidity:  $\eta = \frac{1}{2} \log \frac{z_q}{z_{\bar{q}}}$
- Transverse momentum:  $k_\perp$
- Azimuth:  $\phi$

In the soft-collinear approximation

$$d\mathcal{P} = \frac{\alpha_s(k_\perp) C_F}{\pi^2} d\eta \frac{dk_\perp}{k_\perp} d\phi$$

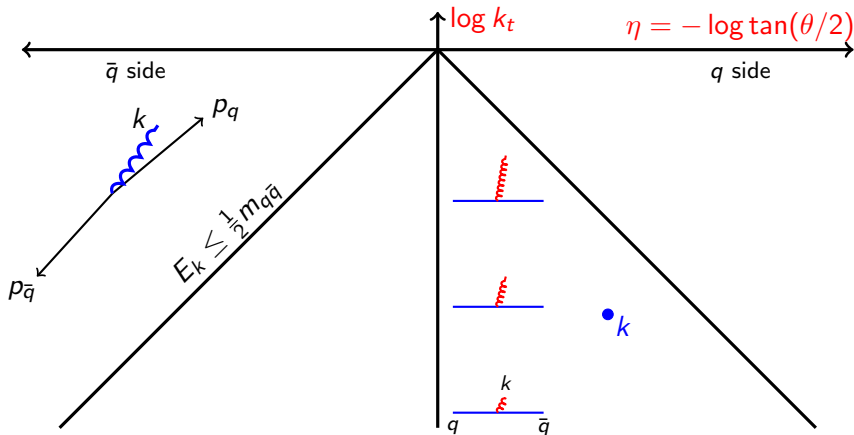
# Basic features of QCD radiations: the Lund plane

Lund plane: natural representation uses the 2 “log” variables  $\eta$  and  $\log k_{\perp}$



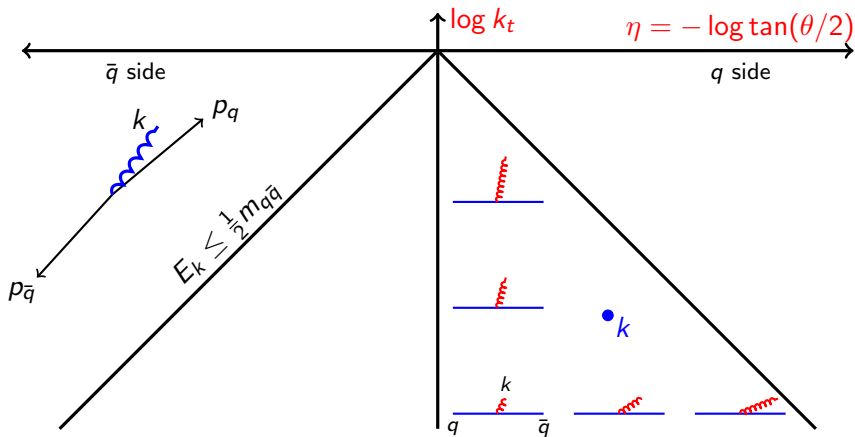
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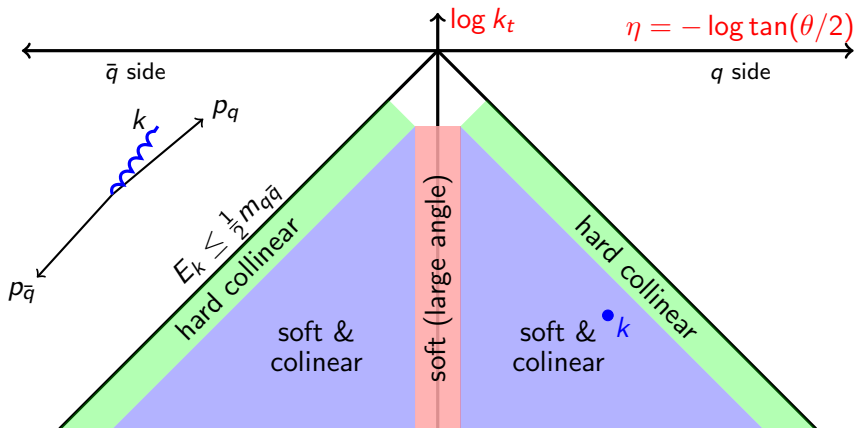
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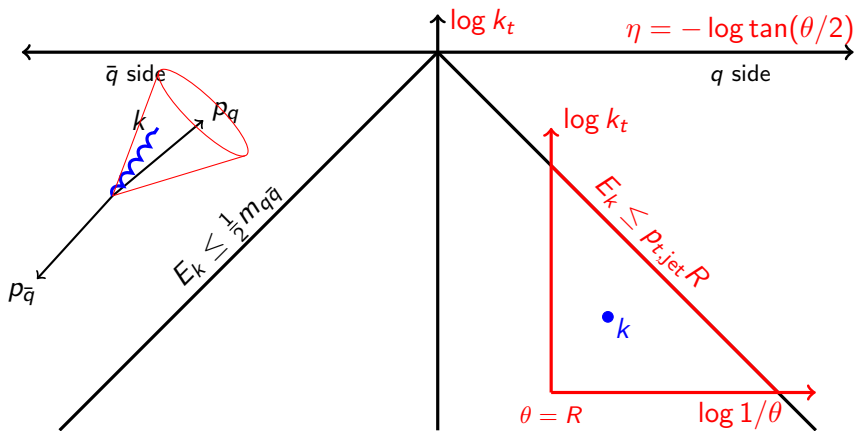
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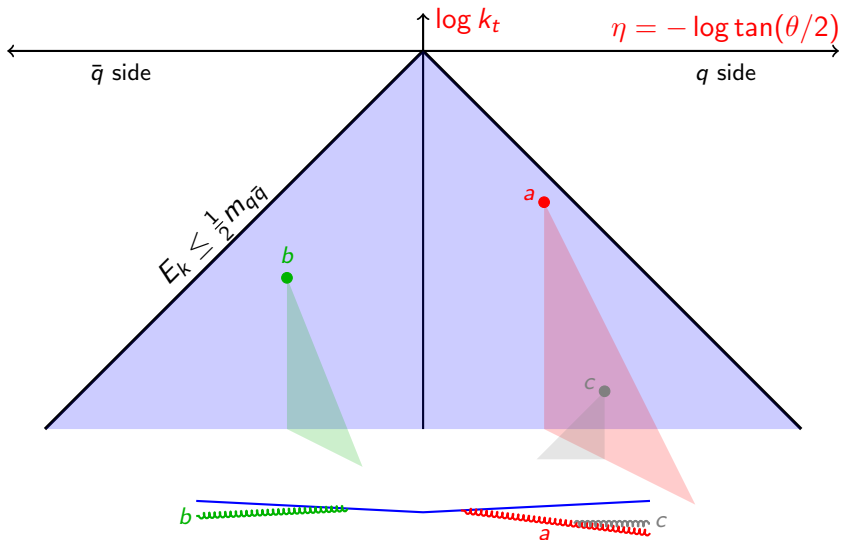
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For a jet:  $\eta = -\log \tan \theta/2 \approx \log 1/\theta$ ,  $\theta < R$

# Multiple emissions in the Lund plane



# Parton showers



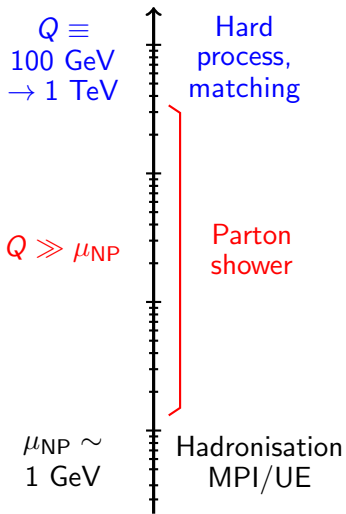
# Parton shower in $pp$ : big picture

Lots of progress over the past decade:

- **1  $\rightarrow$  3 splitting functions**: ingredients towards NLO DGLAP (e.g. Dire)  
See e.g. [Jadach *et al*,16] [Li,Skands,16] [Höche,Krauss,Prestel,17] [Höche,Prestel,17]
- **Beyond leading colour**: most showers (except herwig) use dipoles and are leading  $N_c$  (even at leading log)
  - **Amplitude-level showers** (instead of  $ME^2$ )  
see e.g. [Forshaw,Holguin,Plätzer,19]
  - **Beyond leading- $N_c$ /full colour**  
see e.g. [Nagy,Soper,12] [Gieseke,Kirchgaesser,Plätzer,Siodmock,18]  
[Höche,Reichelt,20] [Forshaw,Holguin,Plätzer,20]
- **Electroweak showers**: include  $W/Z/\gamma$  in showers  
involved splitting functions, explicit dependence on chirality/spin<sup>(\*)</sup>  
see e.g. [Kleiss,Verheyen,20] [Bauer,Ferland,Webber,17-18]  
[Bauer,DeJong,Nachman,Provasoli,19]

(\*) Technically, this is also the case for QCD showers

# What does shower accuracy mean?



“Standard” perturbative expansion

$$\alpha_s(Q)f_1(v) + \alpha_s^2(Q)f_2(v) + \alpha_s^3(Q)f_3(v) + \dots$$

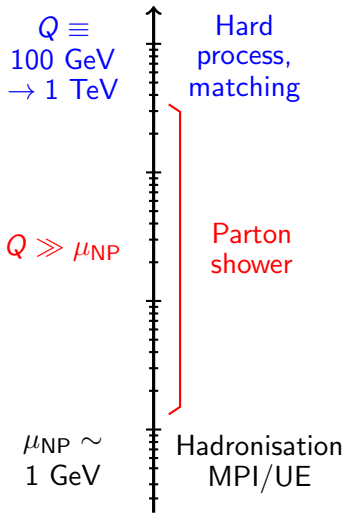
LO                      NLO                      NNLO

expect logs between disparate scales

$$\alpha_s \log^2 Q/\mu_{\text{NP}}, \alpha_s \log Q/\mu_{\text{NP}}$$

(double, single,...) logs to resum

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$$\alpha_s \log^2 Q/\mu_{\text{NP}}, \alpha_s \log Q/\mu_{\text{NP}}$$

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**PS accuracy means logarithmic accuracy: LL, NLL, N<sup>2</sup>LL, ...**

**well-defined**

**+ systematically improvable**

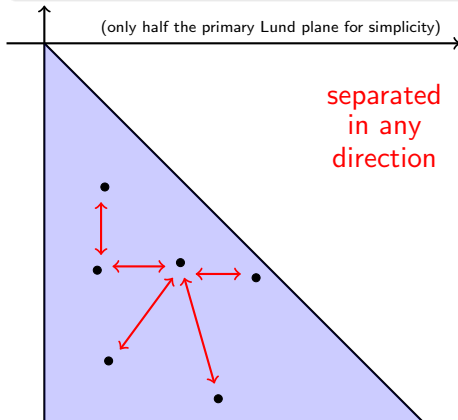
## NLL accuracy for a range of observables

- global event shapes
  - ▶ thrust
  - ▶ jet rates
  - ▶ angularities
  - ▶ broadening
  - ▶ ...
- non-global observables
  - e.g. energy in slice
- multiplicity
  - (NLL is  $\alpha_s^n L^{2n-1}$ )

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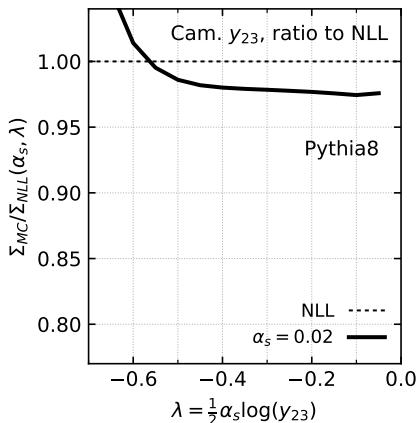
## Correct matrix elements for $N$ well separated emissions in the Lund plane



# Target NLL accuracy

(Cumulative) distributions can (often) be written as

$$P(v < e^{-L}) = \exp \left[ \underbrace{g_1(\alpha_s L)L}_{\text{leading log(LL)}} + \underbrace{g_2(\alpha_s L)}_{\text{next-to-leading log(NLL)}} + \underbrace{g_3(\alpha_s L)\alpha_s}_{\text{NNLL}} + \dots \right]$$



Idea for testing:

$$\frac{\Sigma_{MC}(\lambda = \alpha_s L, \alpha_s)}{\Sigma_{NLL}(\lambda = \alpha_s L, \alpha_s)} \quad \text{v.} \quad 1$$

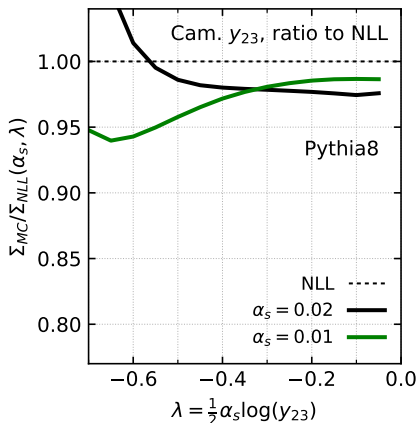
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NLL deviations  
or  
subleading effects?

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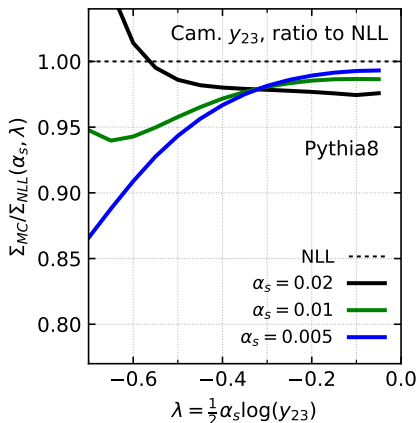
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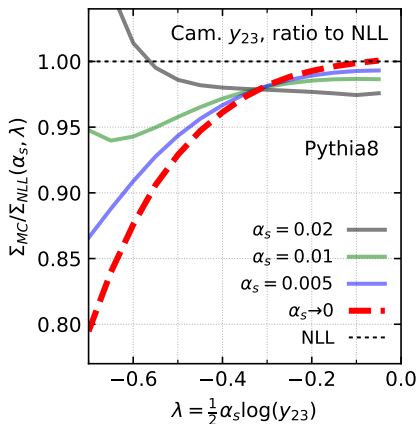
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at fixed  $\lambda = \alpha_s L$

NLL deviations

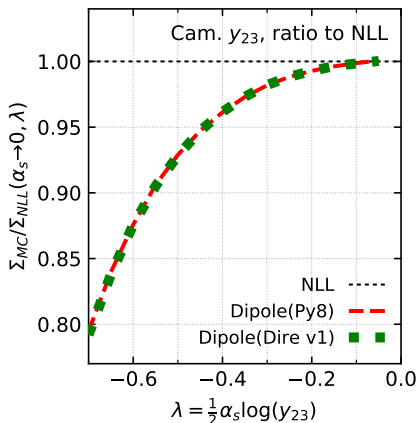
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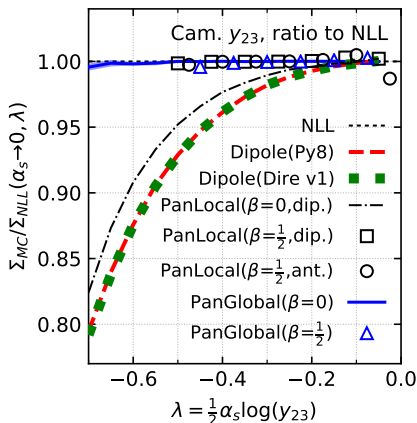
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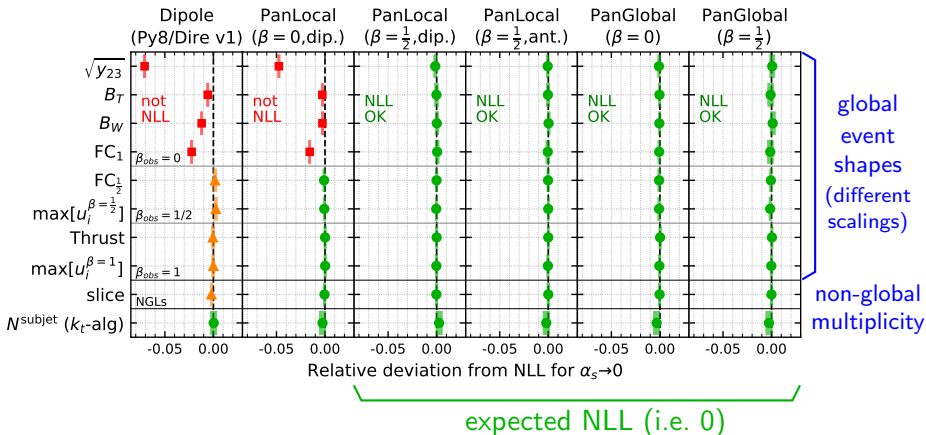
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- × Pythia8 deviates from NLL
- × Dire(v1) same as Pythia8
- ✓ PanLocal( $0 < \beta < 1$ ) OK
- ✓ PanGlobal( $0 \leq \beta < 1$ ) OK

[M.Dasgupta, F.Dreyer, K.Hamilton, P.Monni, G.Salam, GS,20]

# More extensive tests

[M.Dasgupta,F.Dreyer,K.Hamilton,P.Monni,G.Salam,GS,20]



(green: OK at NLL; orange: issues at fixed order; red issues at fixed and all orders)

Recently: prescription beyond leading  $N_c$  [K.Hamilton,R.Medves,G.Salam,L.Scyboz,GS,20]

# Parton shower in HI: big picture

Complex but a good fraction is accessible from first-principles QCD

Significantly improved picture of jet quenching over the past few years

- More precise calculations of medium-induced emissions  
(longitudinal and transverse spectra)
- Accumulate evidence for more fine-tuned effects  
(decoherence, back-reaction, medium response, ...)

see e.g. [Mehtar-Tani,Tywomiuk,19] [Barata,Mehtar-Tani,Soto-Ontoso,Tywomiuk,20]  
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What to look forwards to?

- Still a lot to do “analytically”
  - going beyond simplifying assumptions → higher accuracy/precision
  - more realistic medium description (expansion, geometry, ...)
- More implementations in dedicated HI Monte Carlo generators
- **Benefit from work in generators in  $pp$  collisions**

How to combine (angular-ordered) “vacuum” splittings with (not-angular-ordered) medium-induced emissions?

Idea: compare the transverse momenta over the formation time:  $t_f = \frac{2}{\omega\theta^2}$

$$k_{\perp,\text{vac}}^2 = \omega^2\theta^2$$

$$k_{\perp,\text{med}}^2 = \hat{q}t_f = \frac{2\hat{q}}{\omega\theta^2}$$

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At leading log, 2 possible cases:

- $k_{\perp,\text{vac}}^2 \gg k_{\perp,\text{med}}^2$ : standard vacuum emissions (angular-ordered)(\*)
- $k_{\perp,\text{vac}}^2 \ll k_{\perp,\text{med}}^2$ : medium-induced emission

transition at  $k_{\perp,\text{med}}^2 = k_{\perp,\text{vac}}^2$  i.e.  $\omega^3\theta^4 = 2\hat{q}$

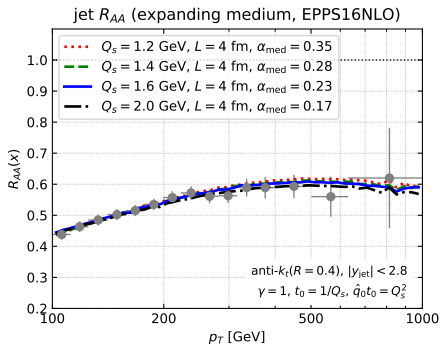
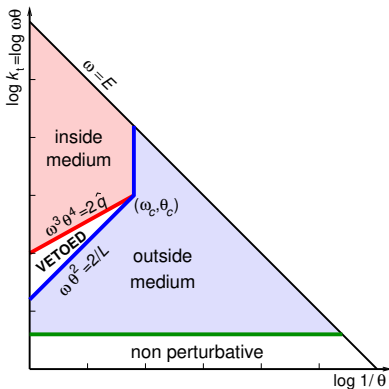
[P.Caucal,E.Iancu,A.H.Mueller,GS,17] (\*) includes details about colour coherence)



# Factorised physical picture

3-stage picture (at double-log):

- 1 in-medium angular-ordered vacuum emissions
- 2 each parton sources MIEs propagating through the medium
- 3 out-medium VLEs with first emission at any angle



includes nuclear PDFs and expanding medium

## A few ideas to keep in sight for future developments

- Helpful to think in terms of ordered scales

$$Q \gg \sqrt{\hat{q}L} \gg T$$

to factorise different effects in a systematic way

- Can one devise accuracy tests similar to the ones discussed for  $pp$ ?

# Jet substructure

# Substructure opens (almost) endless options

## Brief history

1980	Birth
2008	Re-birth (BDRS)
2008-13	Main techniques
2013	First analytics
2013-	New techniques
2018	Deep-learning
2018	Heavy-ions

## Main interest

Offers a differential view of  
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## What existing techniques are good for

Study the dynamics of the QCD branchings in many different ways

Caveat: substructure tools affect quenching effects in non-trivial ways

## Where existing techniques are limited

Jet quenching effects are different from  $pp$  parton shower: angular-ordering violations, different phase-space, ...

Caveat: delicate to find observables which isolate a given quenching effect

# Substructure example: radiation in the Lund plane

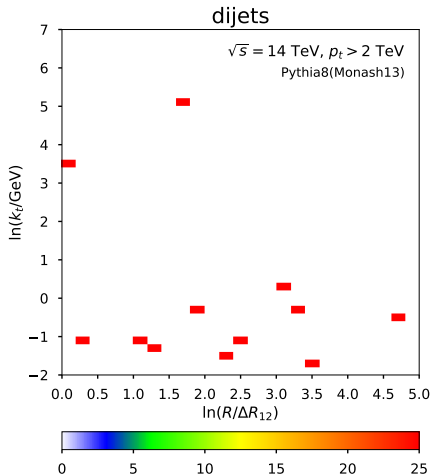
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## The (primary) Lund jet plane

- Cluster the jet with Cambridge/Aachen (i.e. orderd in angles)
- Iteratively undo last clustering (following hardest subjet)
- measure  $k_t$  and  $\Delta R$  of branching



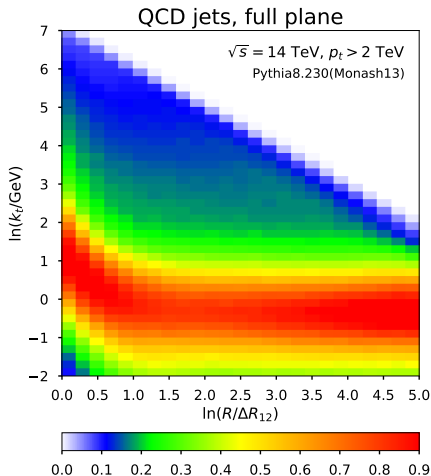
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One (of many) nice properties:  
separate different physical regions



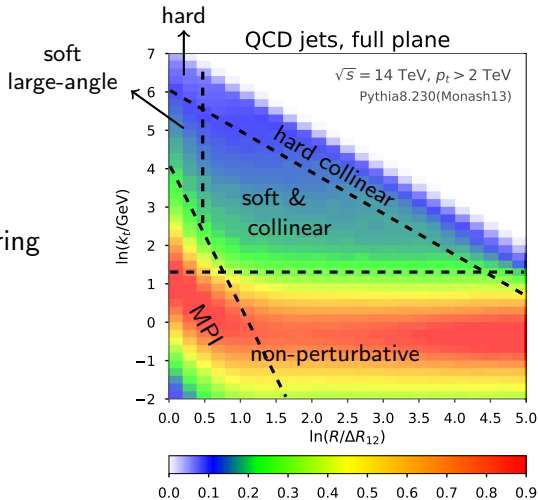


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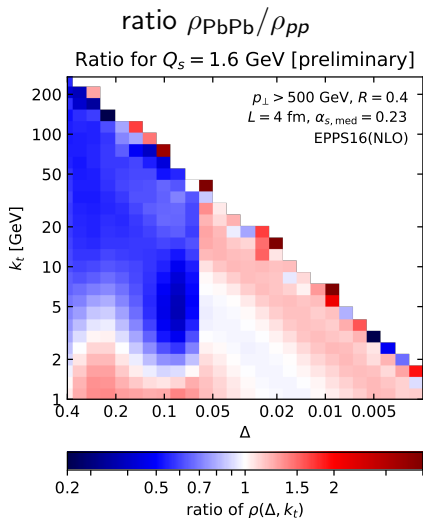
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# Lund plane in heavy ions

Lund plane potentially interesting in heavy-ion collisions

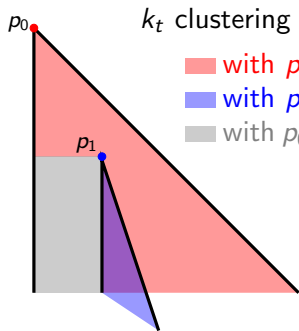
- picture not 100% clear
- some effects
  - quark v. gluon  $E_{\text{loss}}$
  - decoherence
  - in/out interface
  - medium induced
- beyond primary plane?
- correlations?



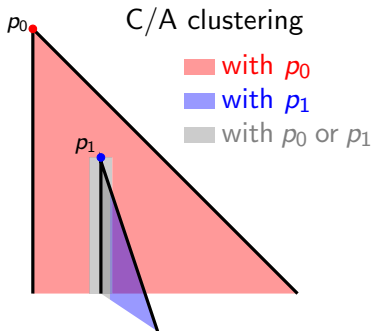
# Substructure: (non) angular ordering and C/A clustering

In  $pp$ , C/A is natural since it respects (strong) angular ordering

Imagine a leading parton  $p_0$  + a largest- $k_t$  emission  $p_1$



“wrong” clustering can happen  
at double-logaccuracy



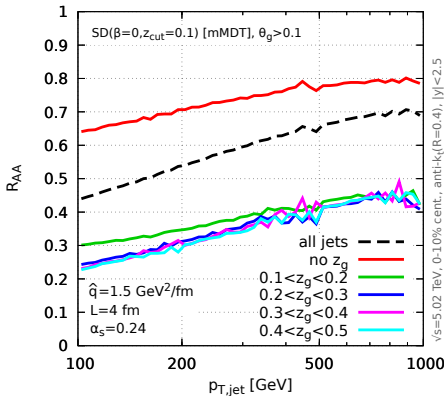
clustering as expected  
except in “single-log” region

## What are the appropriate substructure tools in HI?

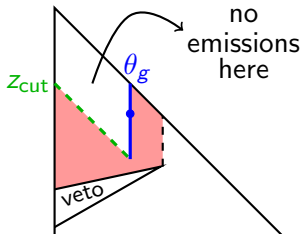
- Large part of the phase-space dominated by vacuum
  - ▶ Can be an advantage: *use substructure to select the vacuum configuration that propagates through the medium*
  - ▶ C/A still appears as a natural choice (for clustering-based tools)
- Different ordering can be probed
  - ▶ C/A declusterings gives an ordered list of declusterings (or a full tree); *these can be ordered as we please* (cf. e.g. dynamical grooming)
  - ▶ the first clustering is unaffected  $\Rightarrow$  different algorithms can be used

## Idea: correlate jet quenching with substructure variables

Dependence of  $R_{AA}$  on  $z_g$  ( $\theta_g > 0.1$ )



This is just an example...



$R_{AA}$  in different bins of  $\theta_g$

- smaller  $\theta_g \Rightarrow$  less sources  
 $\Rightarrow$  less  $E_{\text{loss}}$
- transition at  $\theta_g \sim \theta_c$   
(decoherence angle)

## genuine connections between $pp$ and HI collisions

### Where HI could think $pp$

- large phase-space for “ $pp$  physics”
- be quantitative
  - ▶ assess accuracy
  - ▶ include uncertainties

### Where $pp$ could think HI

- rich QCD pheno of QGP interactions
- interesting challenges to think about e.g. using substructure