

Jets and heavy flavours: an introduction

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- **generic introduction:** QCD and strong interactions
What is QCD? Why do we need it (at the LHC)?
- **heavy flavour:**
Why is it special, important, interesting?
What are the main issues?
- **jets:**
Why is it important? What are the main issues?

Strong interactions

Quantum Chromodynamics: basics

QCD is the quantum theory for strong interactions

	QED	QCD
matter	e, μ, τ	6 quarks flavours u, d, s, c, b, t
vector	photon	gluon
quantum nr	charge	colour
sym. group	U(1)	SU(3)

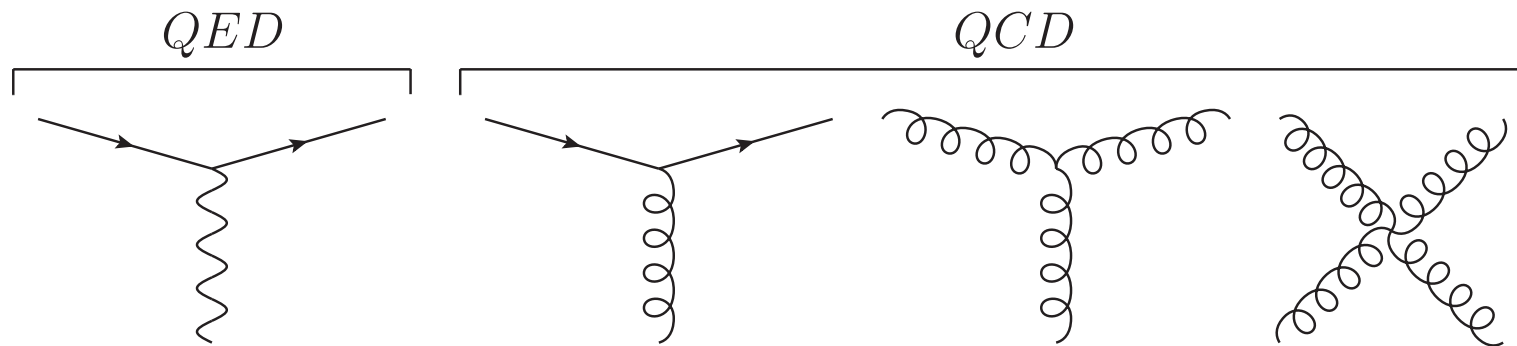
Notes:

- quarks also carry elm charge/interact with photons
- SU(3): 3 fundamental colours (RGB) *i.e.* 3 for quarks, 8 for gluons
- SU(3) is non-abelian

Quantum Chromodynamics: a non-abelian theory

2 main consequences:

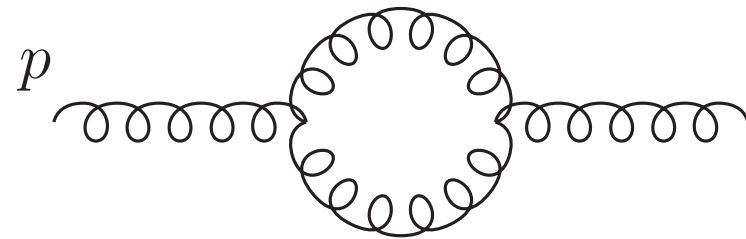
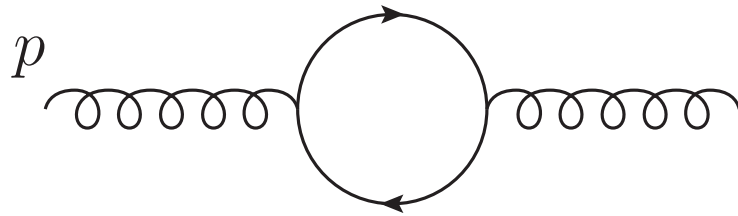
- the gluons interact together



Quantum Chromodynamics: a non-abelian theory

2 main consequences:

- the gluons interact together
- The “running” coupling constant ($\alpha_s = g_s^2/(4\pi)$) decreases with energy



$$\alpha_s(p) = \frac{1}{b_0 \log(p^2/\Lambda_{\text{QCD}}^2)}$$

with

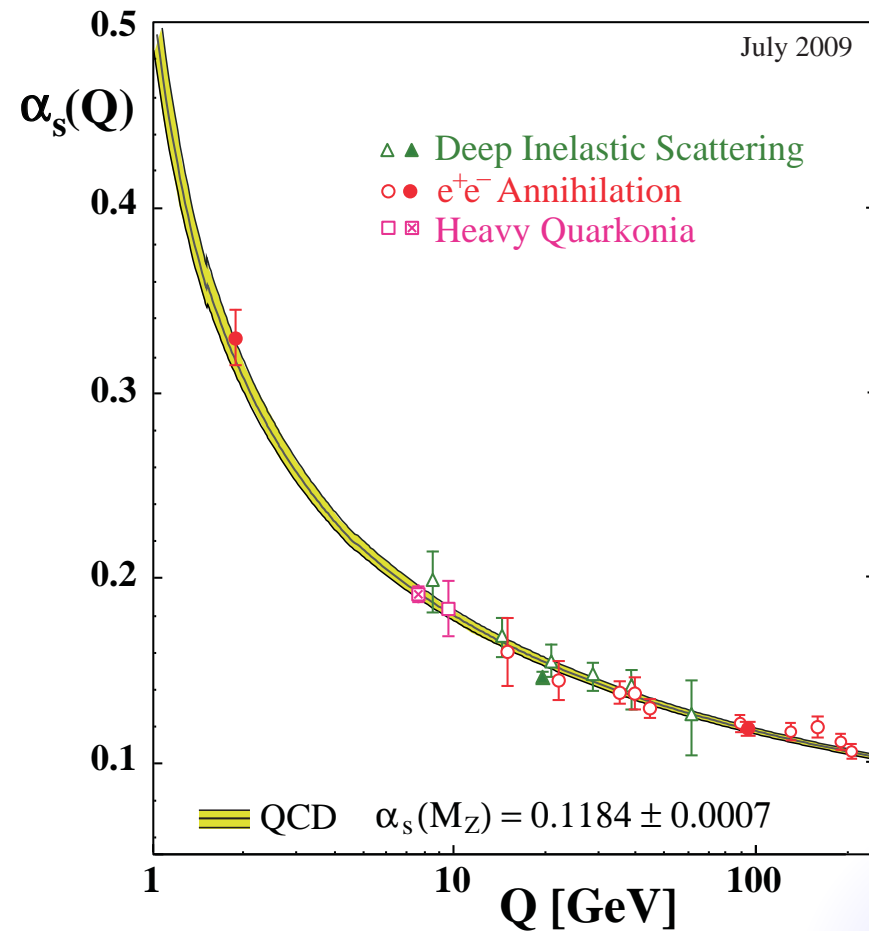
$$b_0 = \frac{11N_c - 2n_f}{12\pi}$$

$b_0 > 0$ for $N_c = 3$ and $n_f = 3 \dots 6$.

Quantum Chromodynamics: a non-abelian theory

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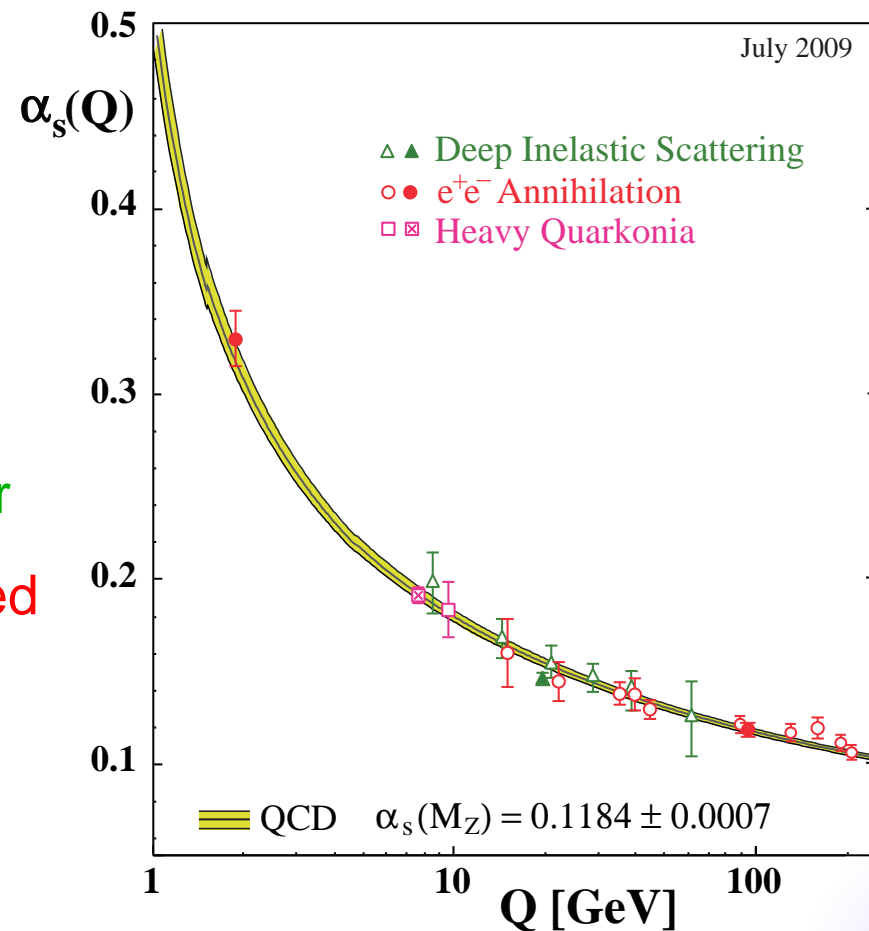
Quantum Chromodynamics: a non-abelian theory

2 main consequences:

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Note:

- $\alpha_s \sim 0.2 \gg \alpha_e$
perturbative corrections larger
- Non-perturbative in the infrared ($\lesssim 1$ GeV)

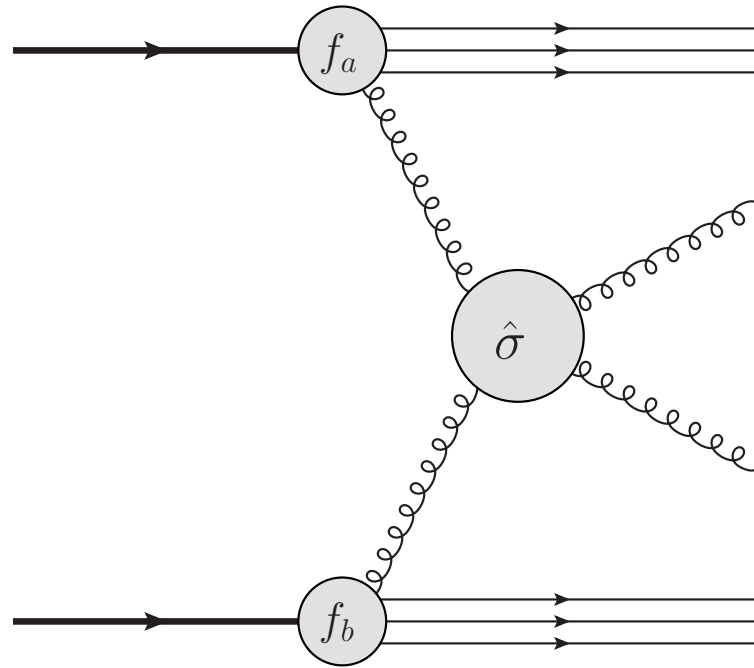


Why is it important at the LHC?

Protons made of quarks and gluons

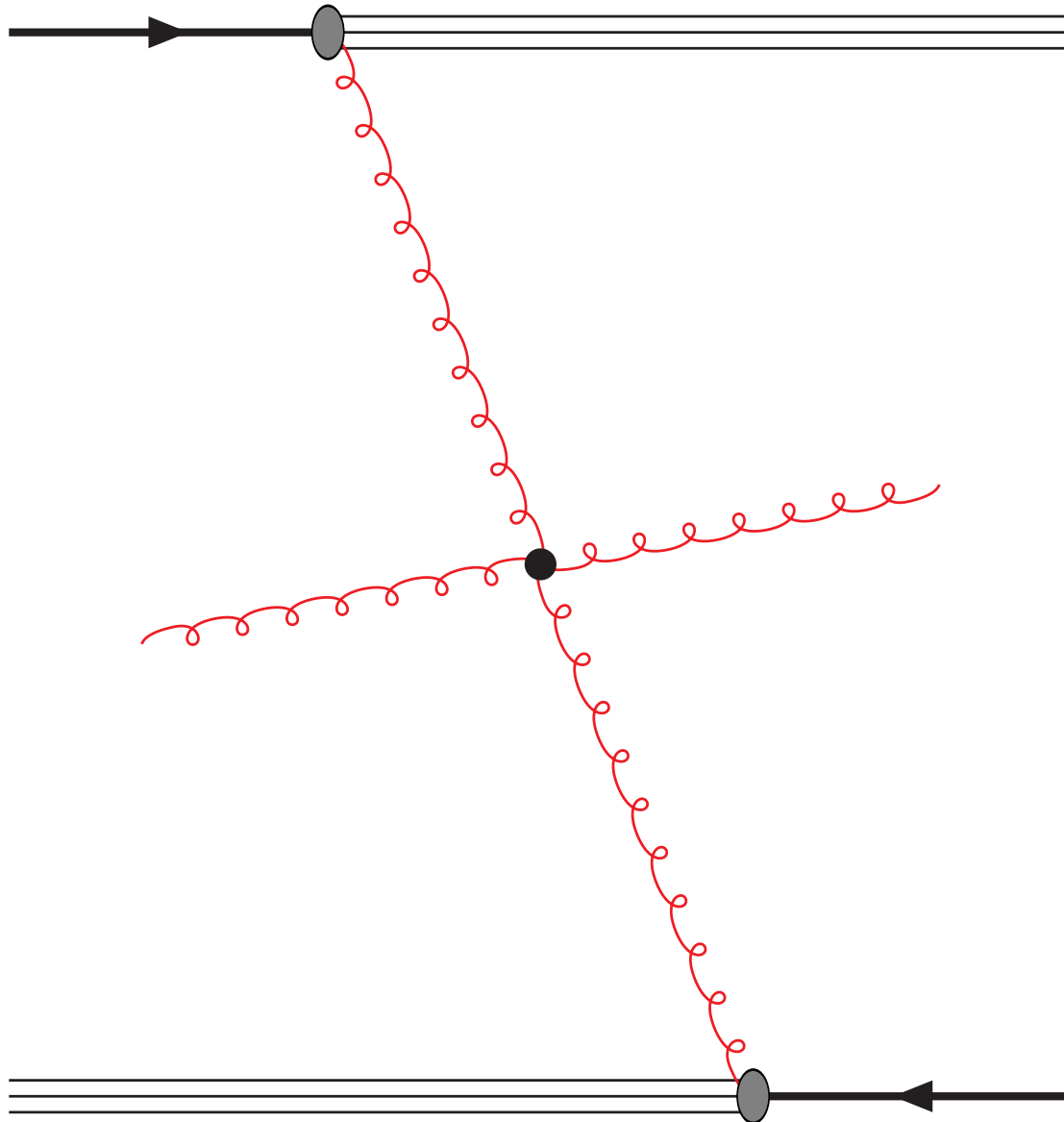
- ⇒ interact mostly through strong interactions
- ⇒ QCD needed for any single event even for electro-weak, Higgs or BSM!

QCD at hadron colliders



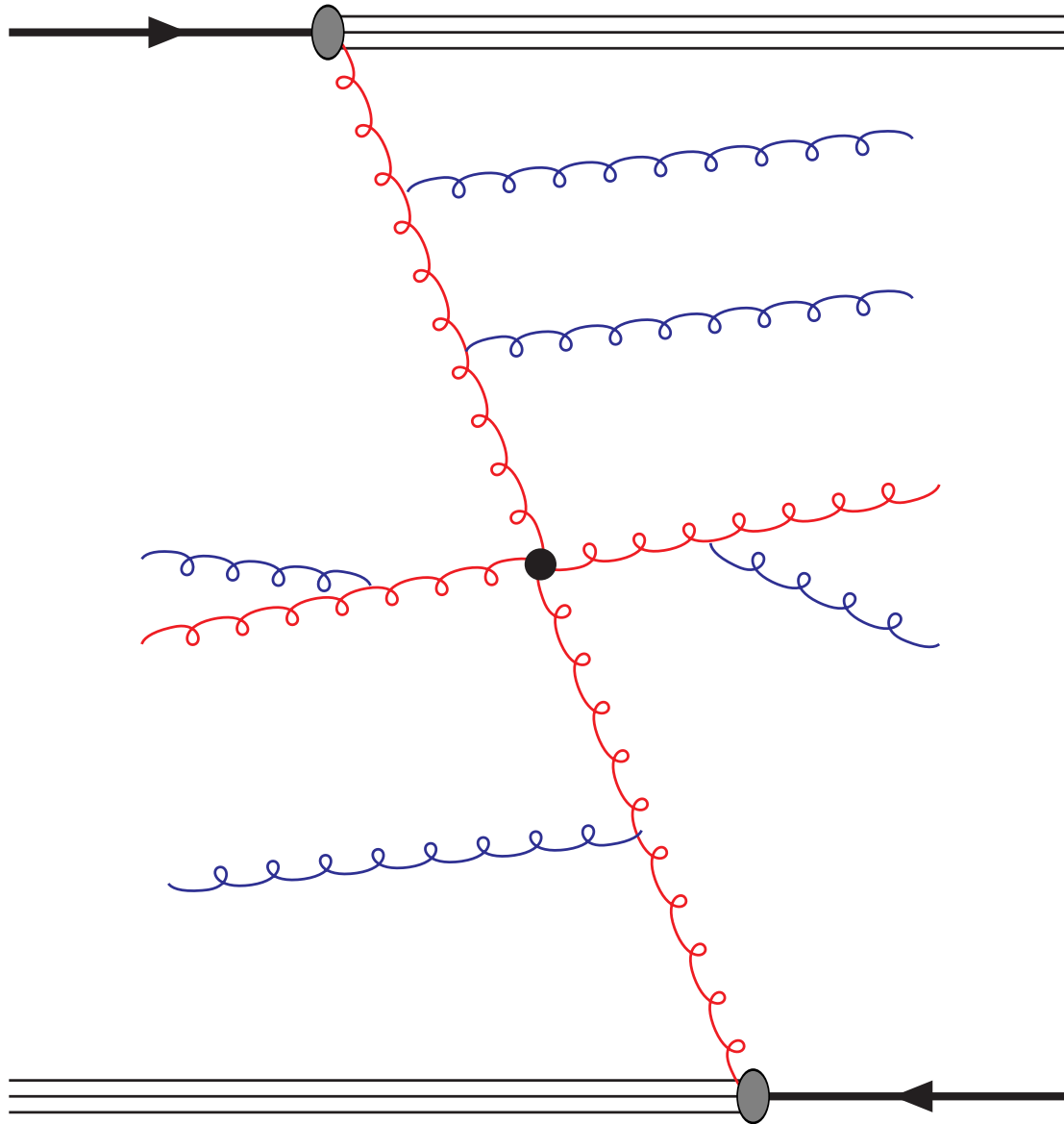
Typical example: QCD needed for the PDF
i.e. the quark and gluon contents of the proton

The more realistic version



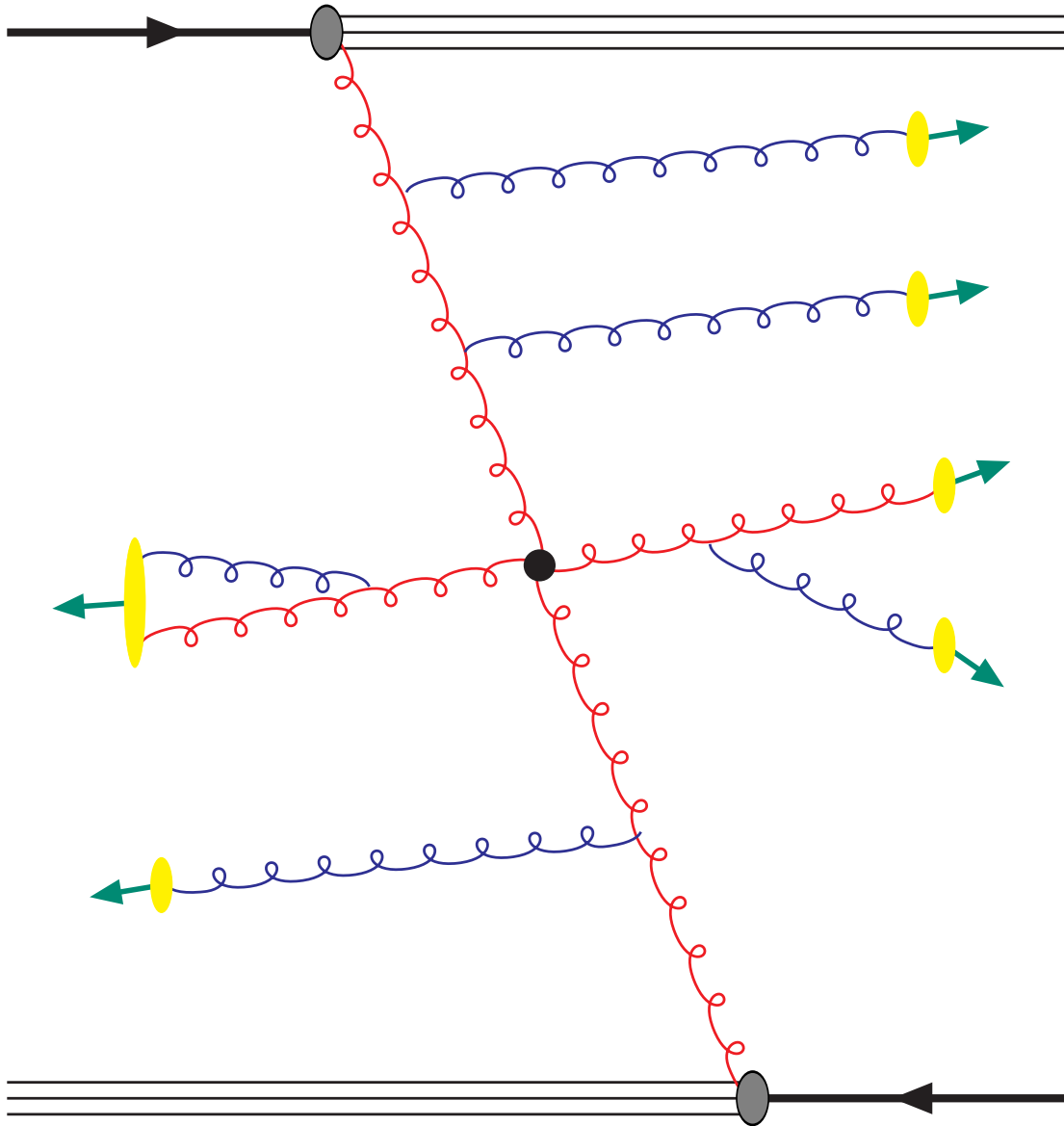
- Hard ME
perturbative

The more realistic version



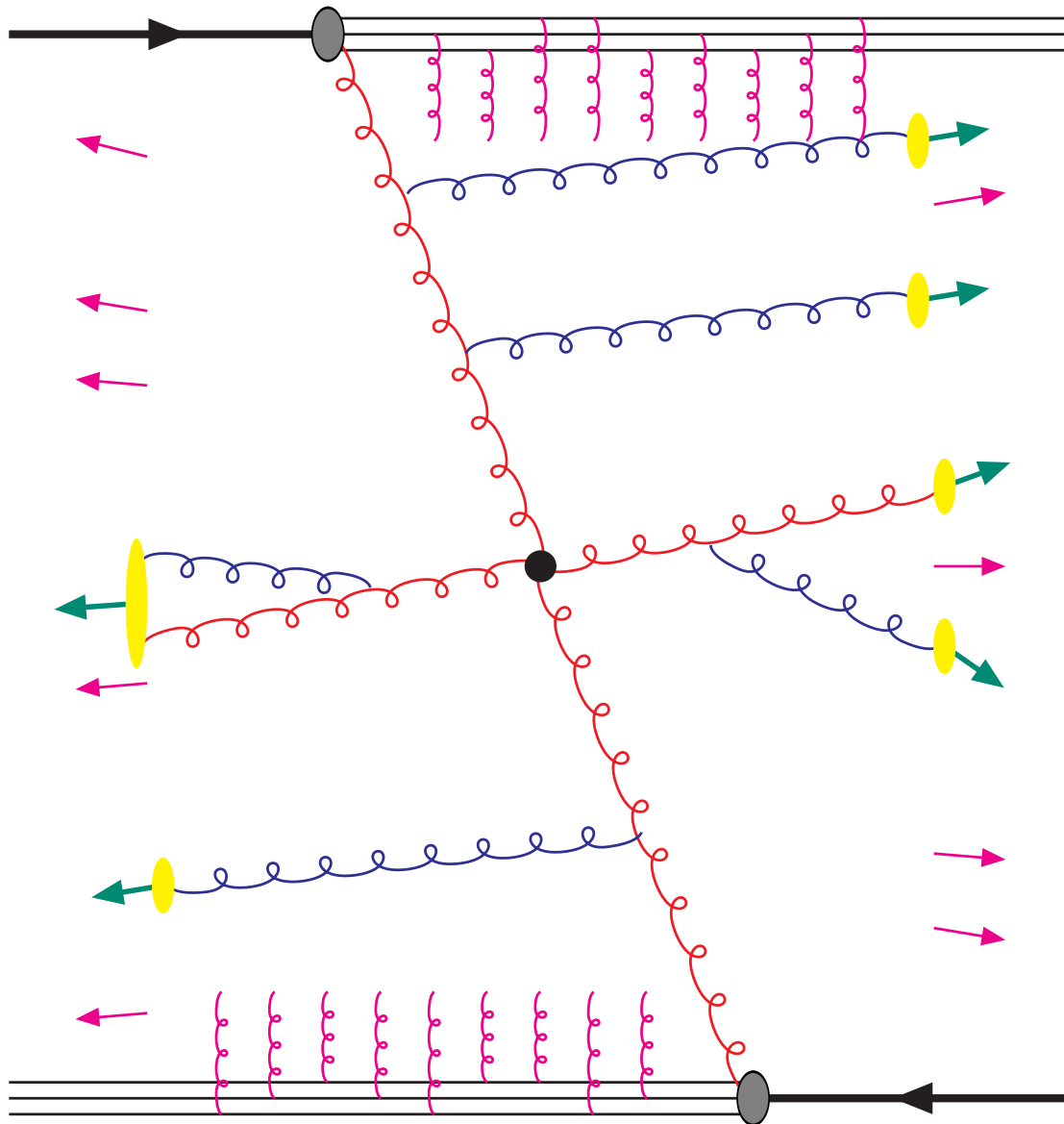
- Hard ME
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- Parton branching
initial+final state radiation

The more realistic version



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 $q, g \rightarrow$ hadrons

The more realistic version



- Hard ME
perturbative
- Parton branching
initial+final state radiation
- Hadronisation
 $q, g \rightarrow$ hadrons
- Multiple interactions
Underlying event (UE)

Heavy quarks

Light and heavy quarks

6 quark flavours:

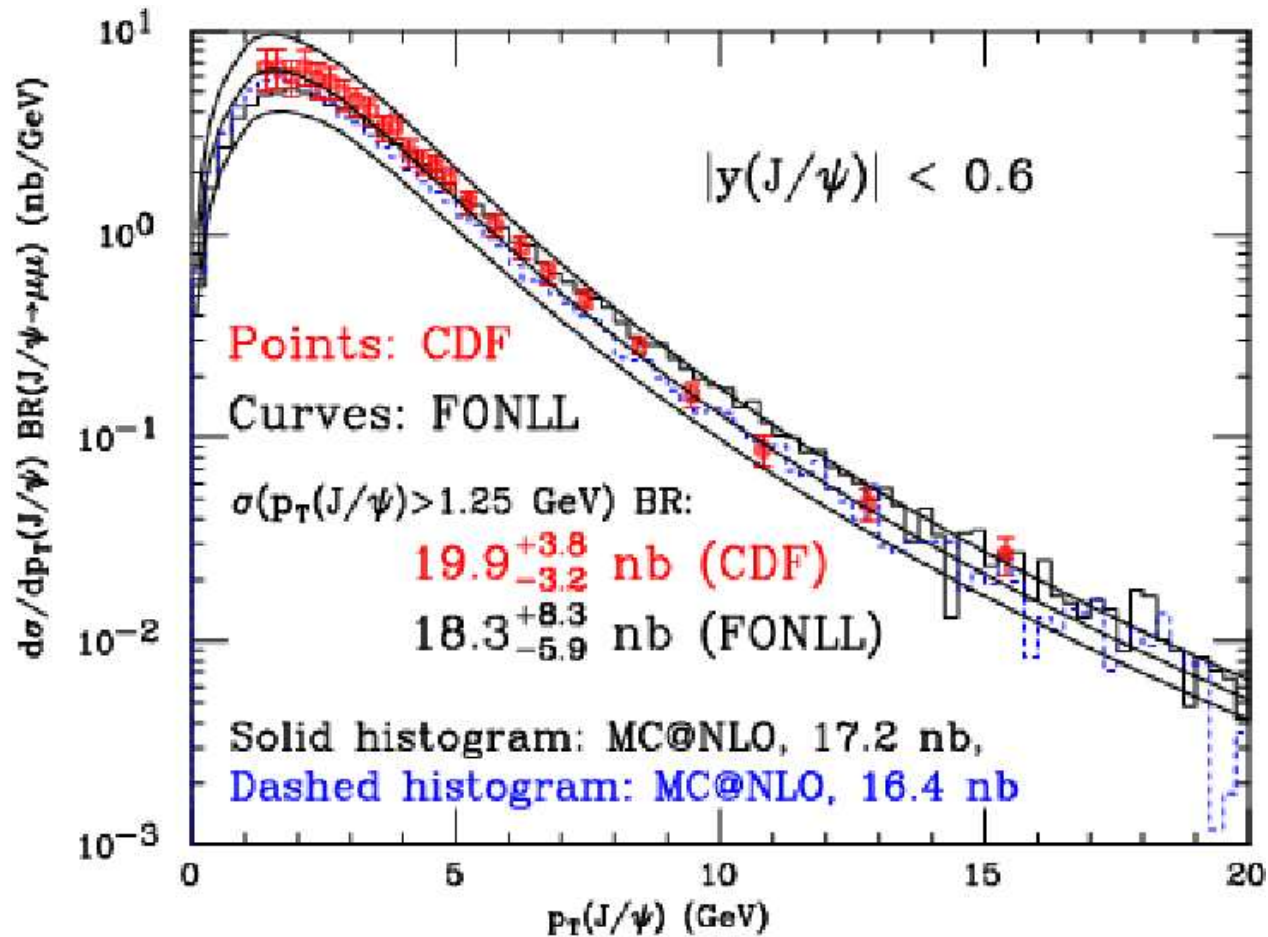
- u, d, s : mass ≈ 0
- c : $m \sim 1.5$ GeV
Etat typique: $J/\Psi \equiv c\bar{c}$
- b : $m \sim 4.5$ GeV
Etats typiques: B, Υ
- t : $m \sim 172$ GeV
Decay into W and b
 $W \rightarrow q\bar{q}$ ($\approx 66\%$), $W \rightarrow \ell\nu$ ($\approx 33\%$)

Why am I here?

- $m_{c,b,t} \gg \Lambda_{\text{QCD}}$: we may apply perturbation theory
- b :
 - could be tagged: displaced vertex
 - SM: b production vs. QCD, top decay
 - new physics search:
 - D0: like-sign $\mu\mu$ charge asym (from b decay)
 - $H \rightarrow b\bar{b}$ dominant at low Higgs mass
- top:
 - top in the standard model *e.g.* mass measurement
 - BSM: coupling $\propto m$
 - \Rightarrow modifications in the top sector
 - \Rightarrow very important at the LHC

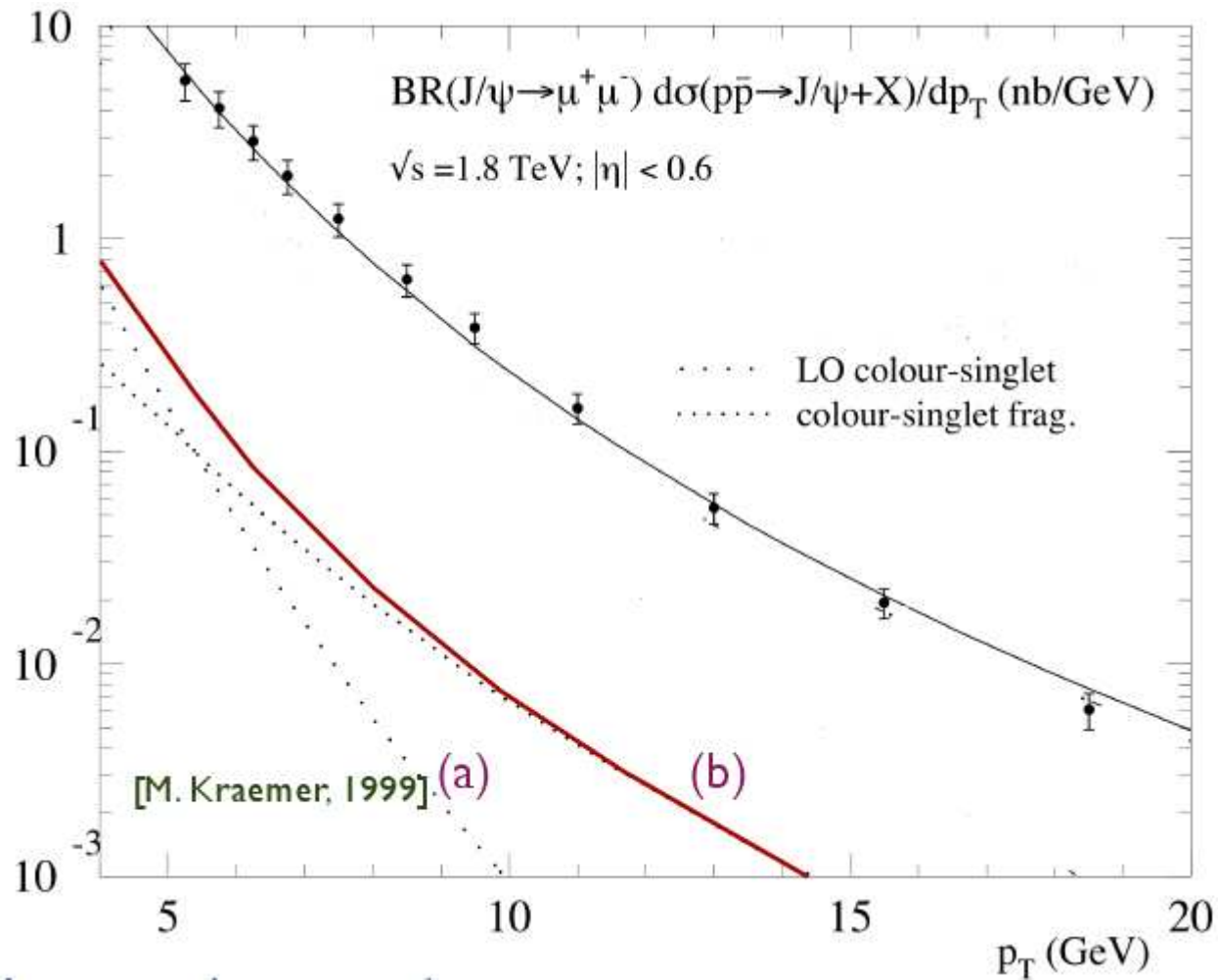
Perturbative QCD: charm

J/Ψ production: $J\Psi$ from b decay



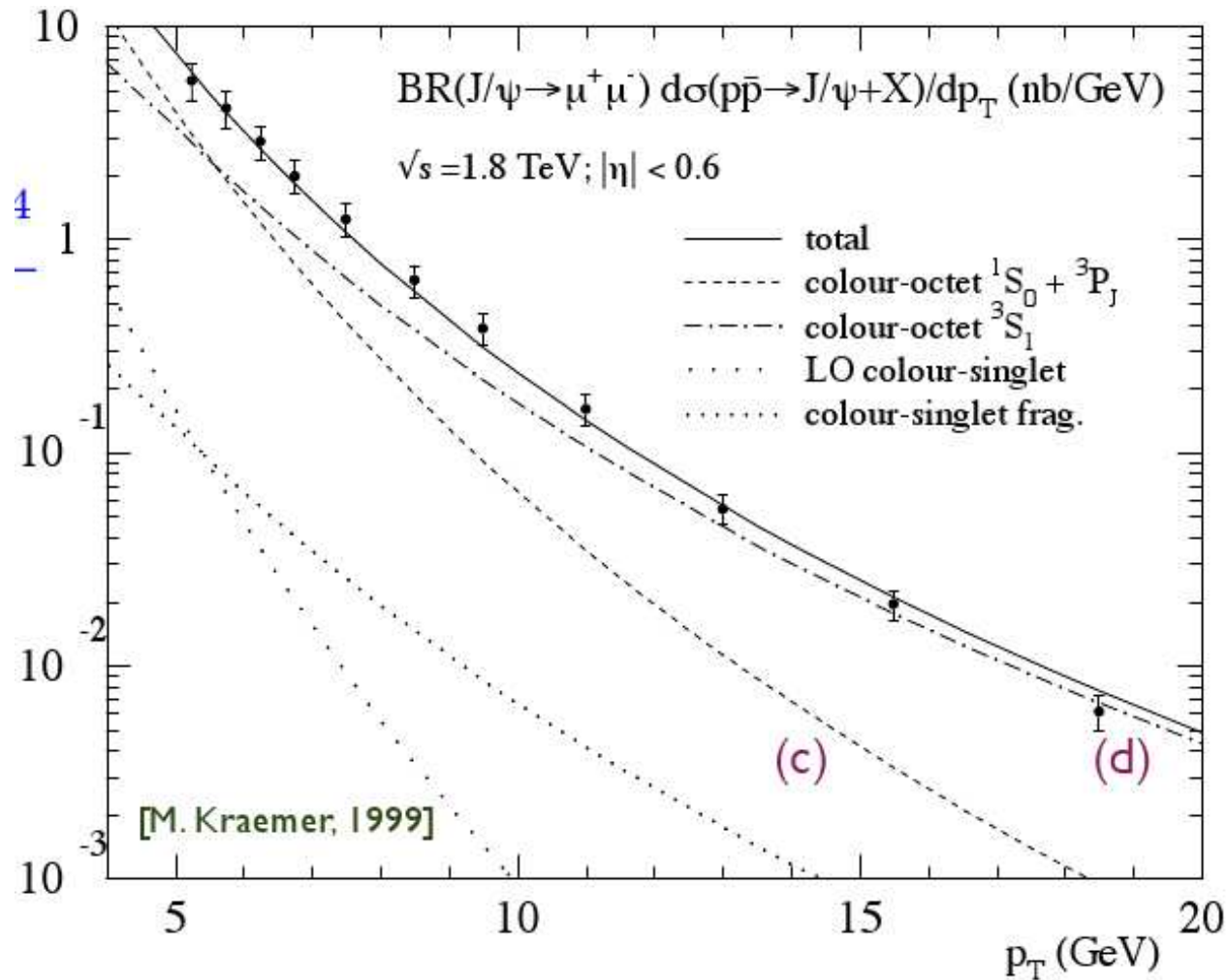
Perturbative QCD: charm

J/Ψ production: **Not the best agreement ever**



Perturbative QCD: charm

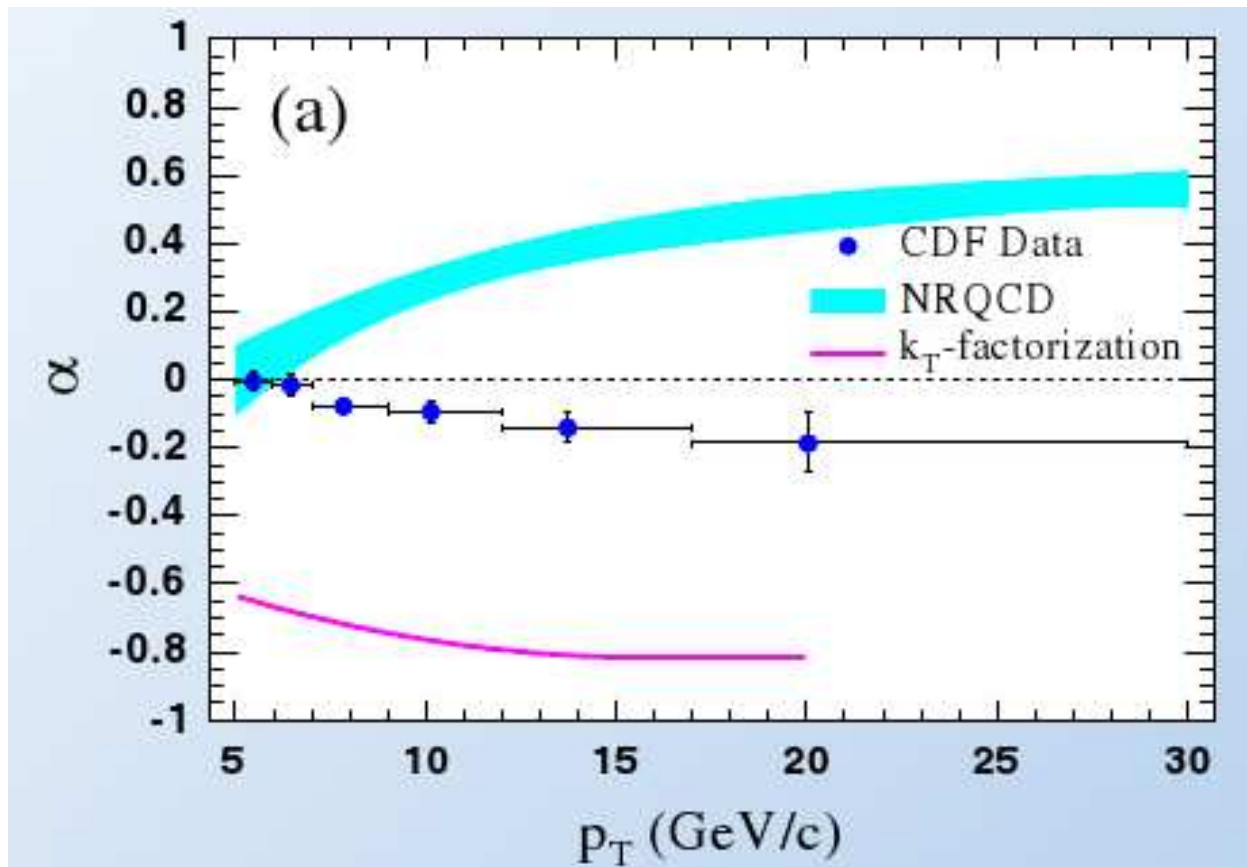
J/Ψ production: better with higher-order corrections*



* agreement not 100% understood

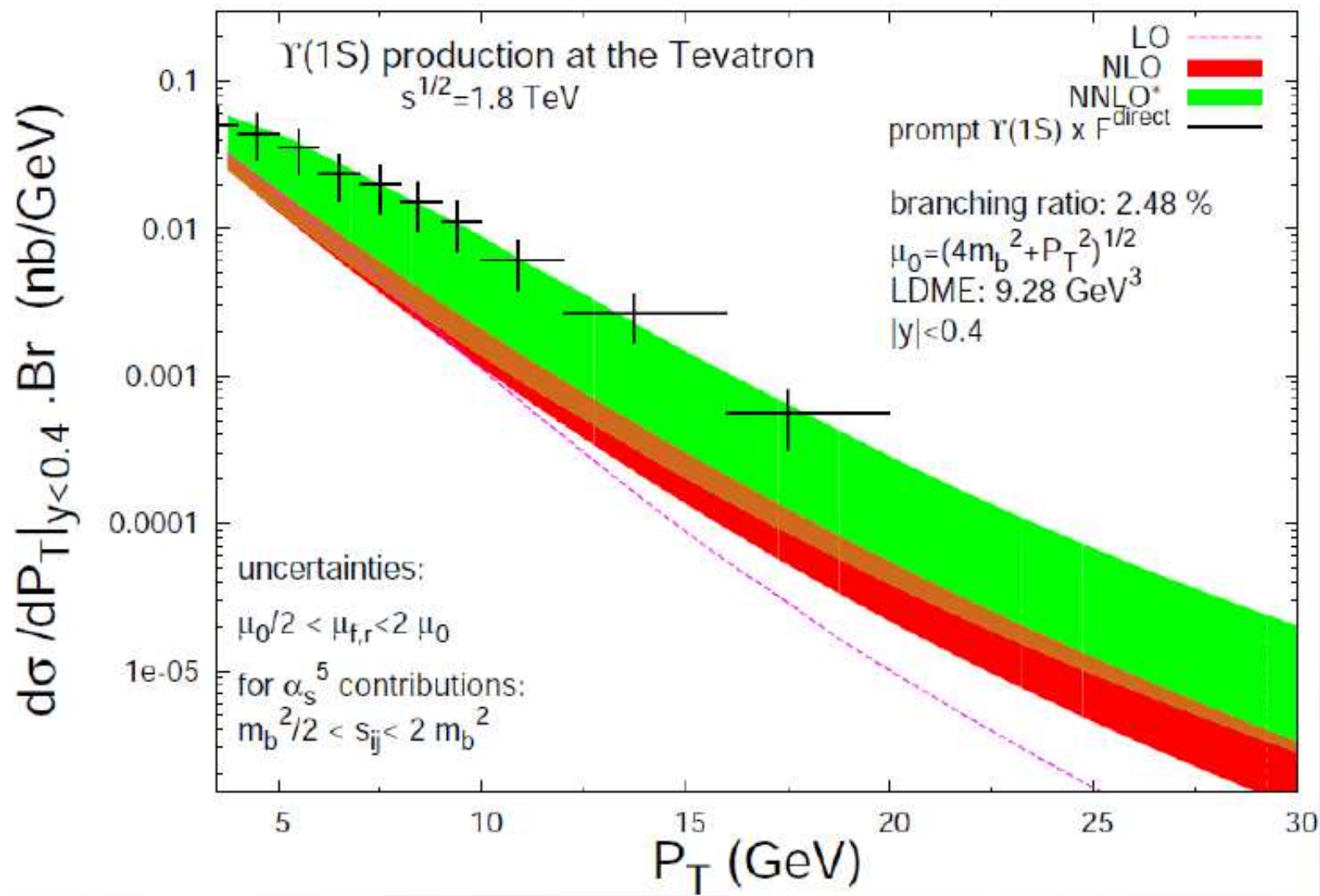
Perturbative QCD: charm

J/Ψ production: room for improvement



Perturbative QCD: bottom

J/Ψ production: again, higher-order important



● Production:

- Mostly $gg \rightarrow t\bar{t}$
- Tevatron: $\sigma_t \approx 10$ pb: **discovery!**
- LHC: $\sigma_t \approx 1$ nb: $\approx 10/s$ **LHC \equiv top factory**

● Decay:

- Mostly $t \rightarrow Wb$
 $t \rightarrow q\bar{q}b$ ($\approx 66\%$) or $t \rightarrow \ell\nu_\ell b$ ($\approx 33\%$)
- for $t\bar{t}$: 3 options
 - **leptonic**: not-so-easy because 2 neutrinos
 - **semi-leptonic**: ℓ , 4 jets (2b) and \cancel{E}_t
(the most convenient)
 - **hadronic**: 6 jets *i.e.* technical to reconstruct
but $\approx 45\%$ of the stat!

“discovery” at the Tevatron

FERMILAB-PUB-94/097-E

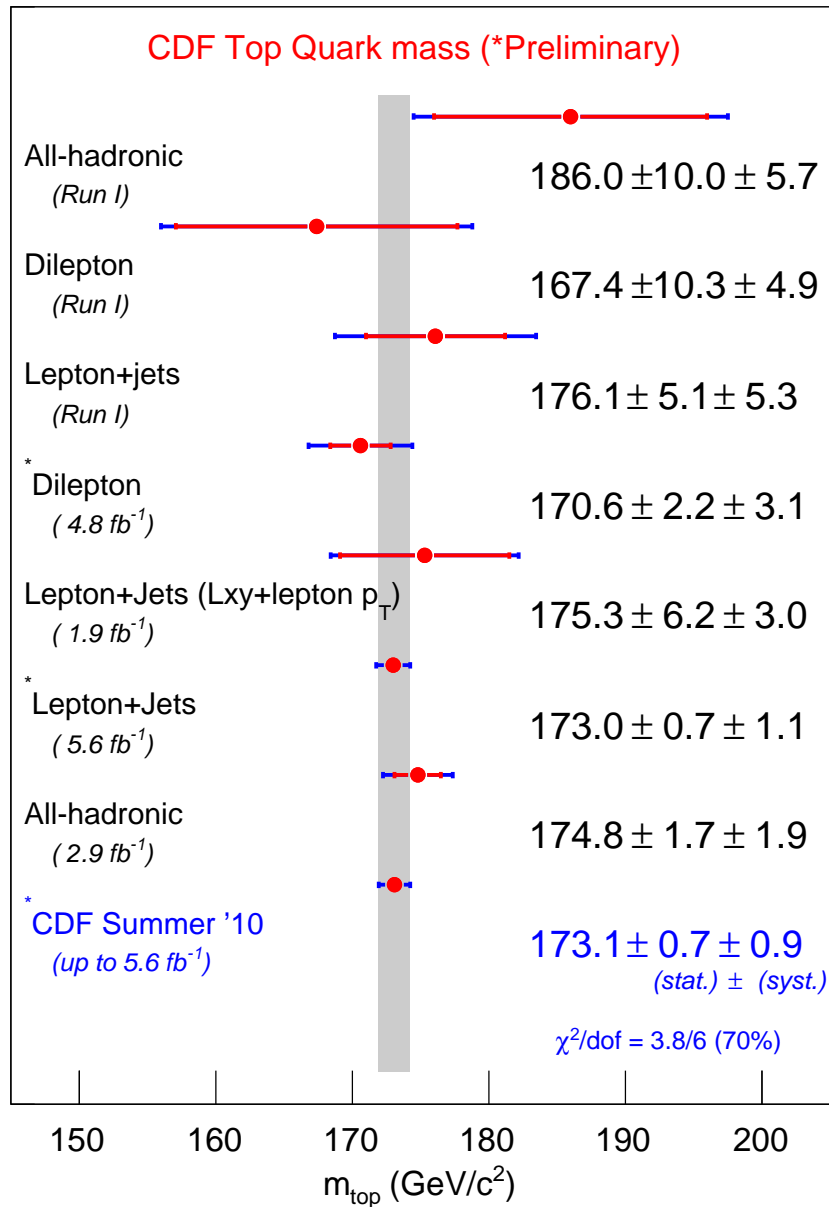
Evidence for Top Quark Production in $\bar{p}p$ Collisions at $\sqrt{s} = 1.8$ TeV

Abstract

We present the results of a search for the top quark in 19.3 pb⁻¹ of $\bar{p}p$ collisions at $\sqrt{s} = 1.8$ TeV. The data were collected at the Fermilab Tevatron collider using the Collider Detector at Fermilab (CDF). The search includes Standard Model $t\bar{t}$ decays to final states $e e \nu \bar{\nu}$, $e \mu \nu \bar{\nu}$, and $\mu \mu \nu \bar{\nu}$ as well as $e + \nu + \text{jets}$ or $\mu + \nu + \text{jets}$. In the $(e, \mu) + \nu + \text{jets}$ channel we search for b quarks from t decays via secondary-vertex identification and via semileptonic decays of the b and cascade c quarks. In the dilepton final states we find two events with a background of $0.56^{+0.25}_{-0.13}$ events. In the $e, \mu + \nu + \text{jets}$ channel with a b identified via a secondary vertex, we find six events with a background of 2.3 ± 0.3 . With a b identified via a semileptonic decay, we find seven events with a background of 3.1 ± 0.3 . The secondary-vertex and semileptonic-decay samples have three events in common. The probability that the observed yield is consistent with the background is estimated to be 0.26%. The statistics are too limited to firmly establish the existence of the top quark, however a natural interpretation of the excess is that it is due to $t\bar{t}$ production. We present several cross checks. Some support this hypothesis, others do not. Under the assumption that the excess yield over background is due to $t\bar{t}$, constrained fitting on a subset of the events yields a mass of $174 \pm 10^{+13}_{-12}$ GeV/c² for the top quark. The $t\bar{t}$ cross section, using this top quark mass to compute the acceptance, is measured to be $13.9^{+6.1}_{-4.8}$ pb.

Top mass today

CDF today



top very important at the LHC

- precision mass measurement
- many new physics scenarios involve the top (mostly because of its large mass)

⇒ need to reconstruct as many tops as possible

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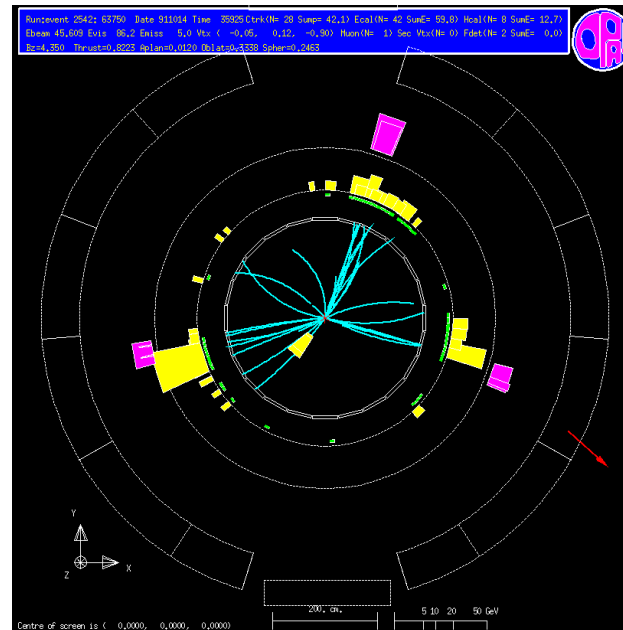
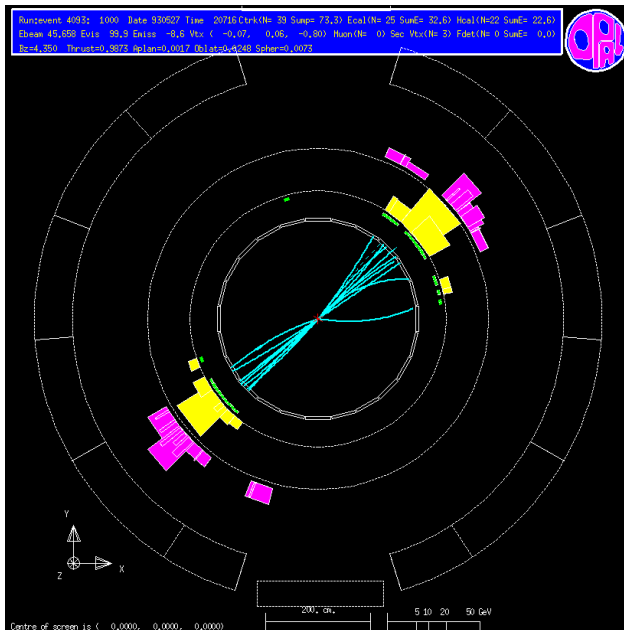
⇒ need to reconstruct as many tops as possible

Issues:

- W +jets background
- b mis-tagging
- combinatorial background (especially for full hadr.)
- efforts e.g. in boosted-top reconstruction

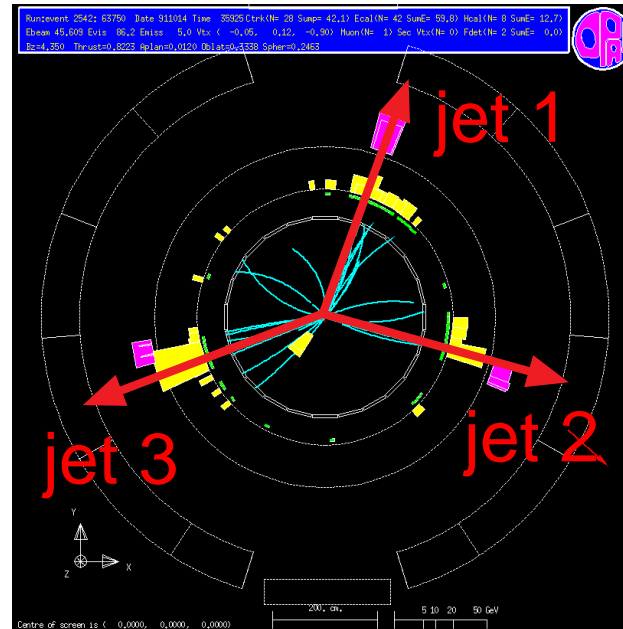
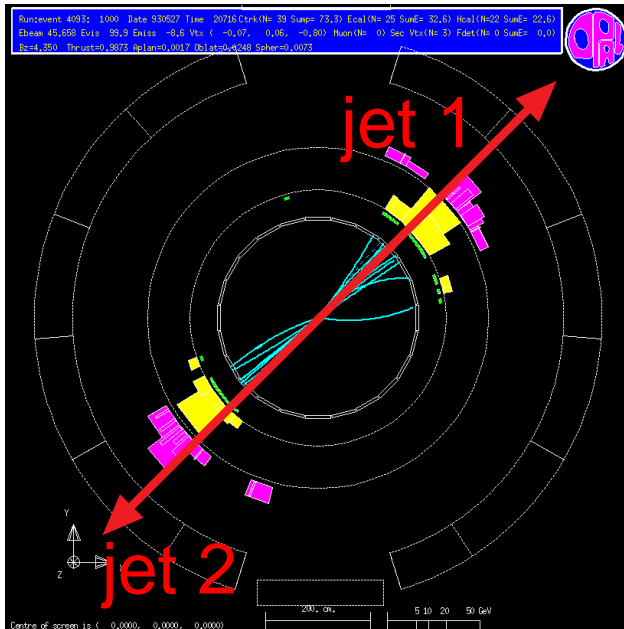
Jets

- Final-state events are pencil-like (already observed in e^+e^- collisions)



- Consequence of the collinear divergence
 QCD branchings are most likely collinear
 $(dP/d\theta \propto \alpha_s/\theta)$

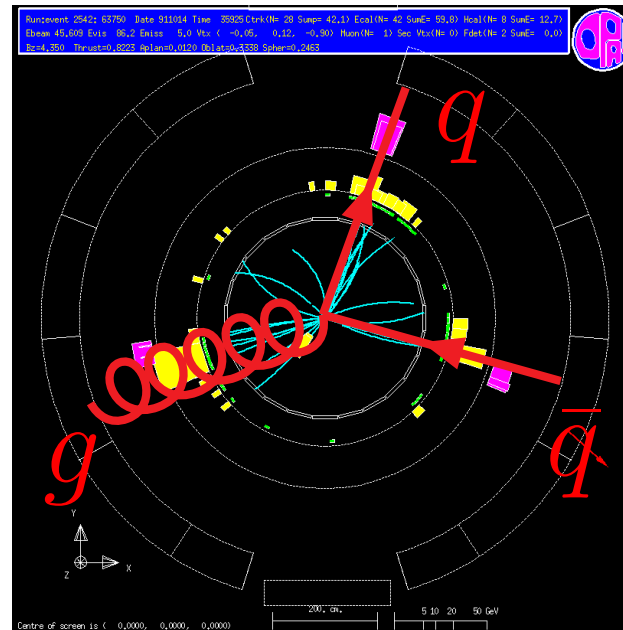
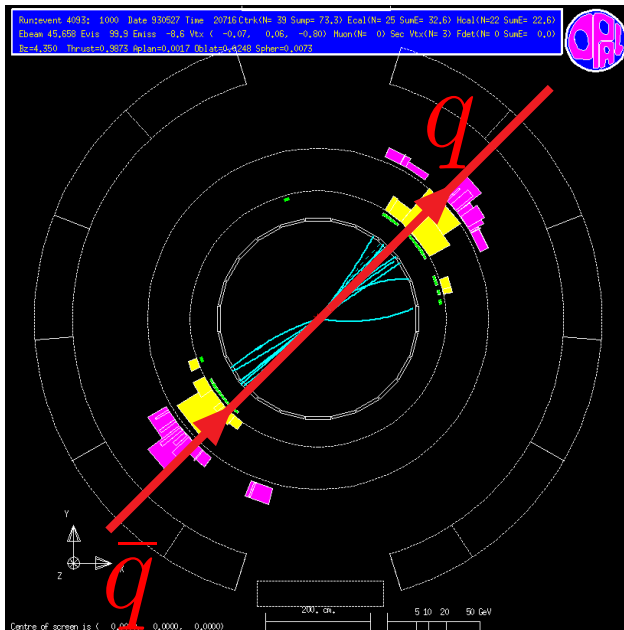
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“Jets” \equiv bunch of collimated particles \cong hard partons

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Why worry

“Jets” \equiv bunch of collimated particles \cong hard partons

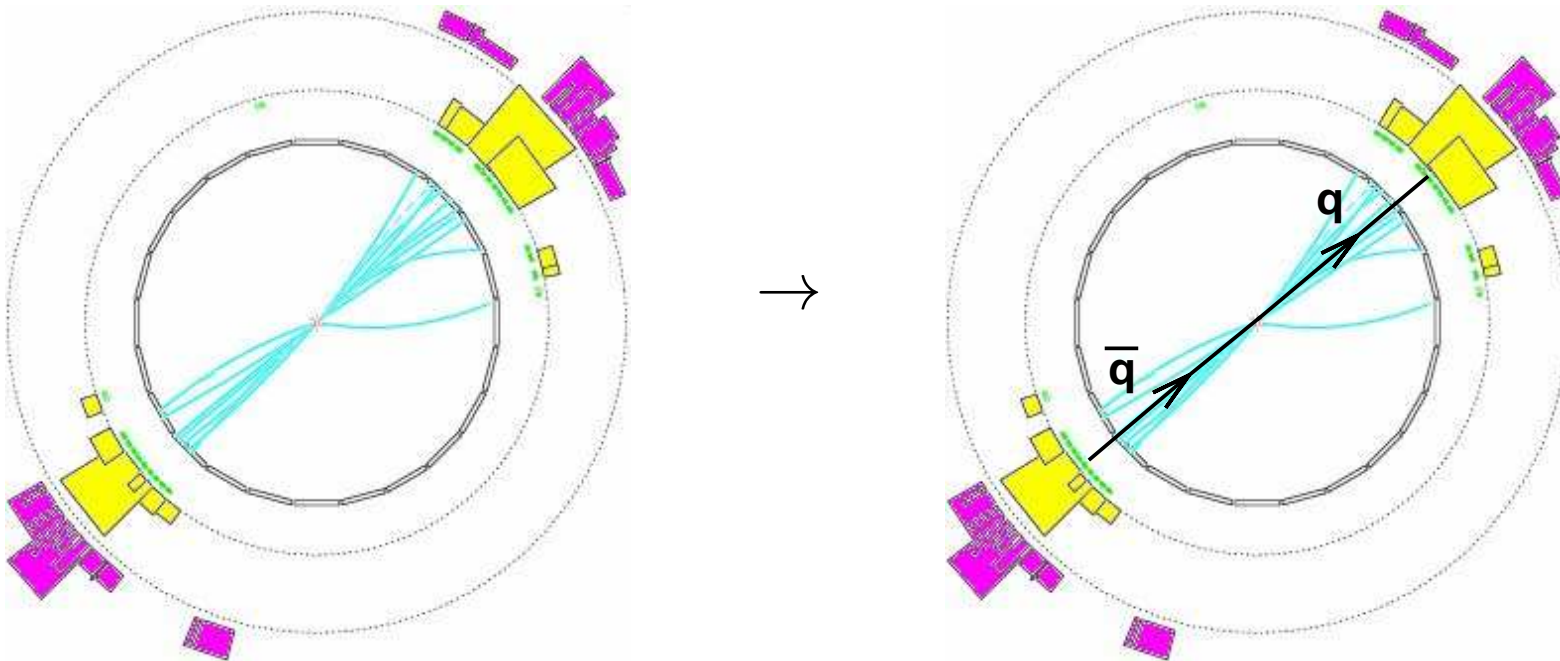
\Rightarrow whenever you have QCD in the final state, you have jets in the final states!

i.e. jets useful mostly everywhere

Jets and partons

“Jets” \equiv bunch of collimated particles \cong hard partons

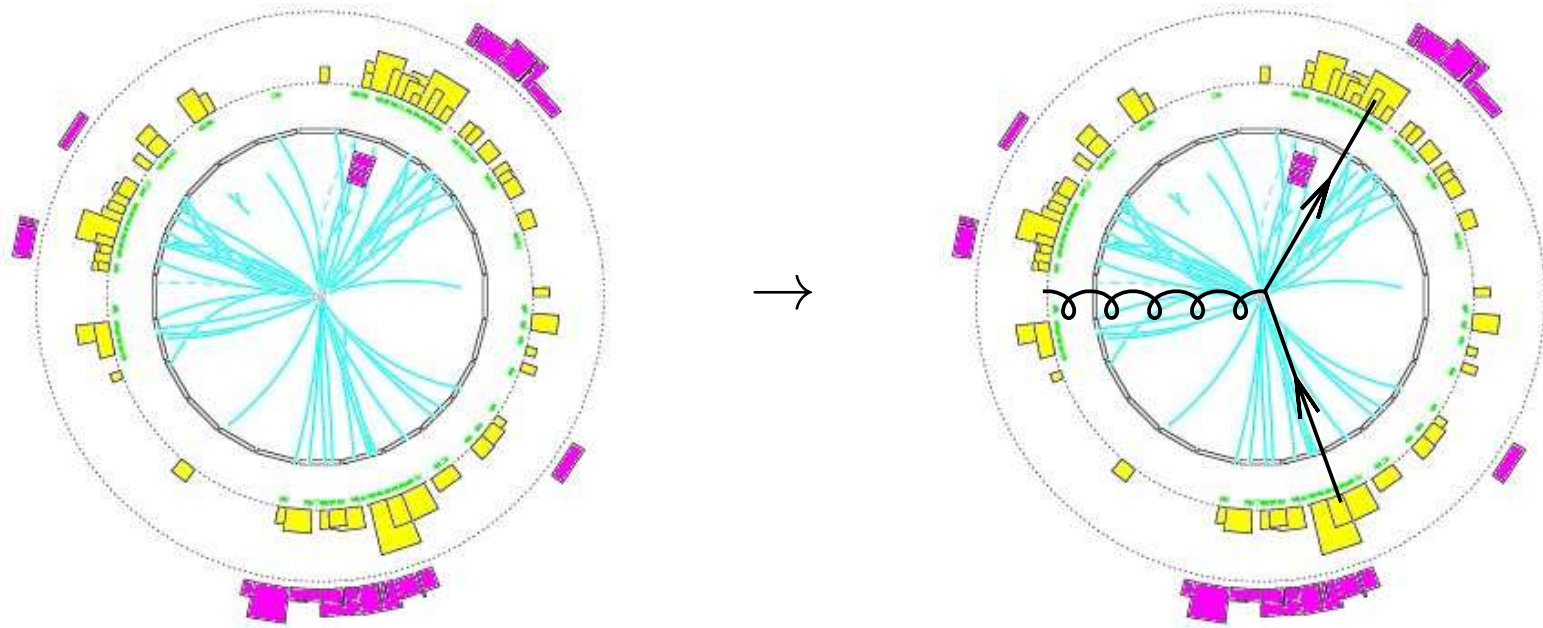
obviously 2 jets



Jets and partons

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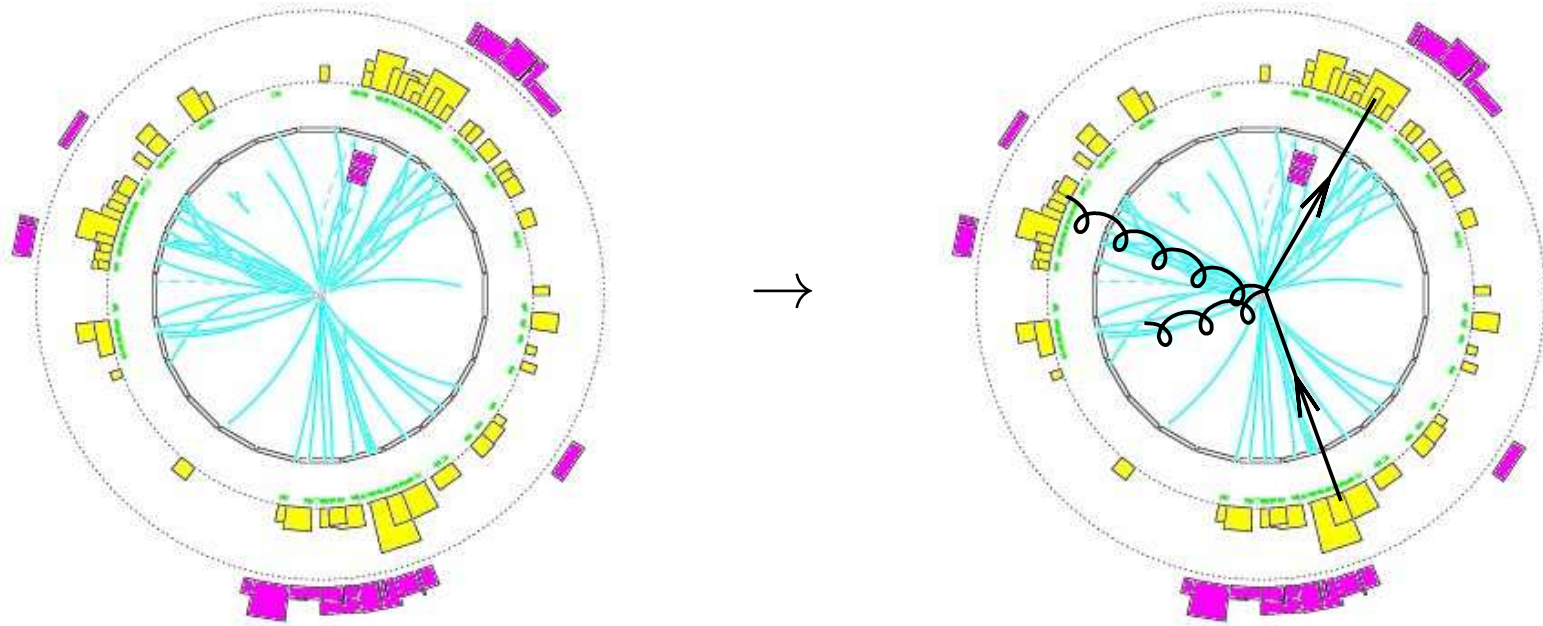
3 jets



Jets and partons

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3 jets... or 4?

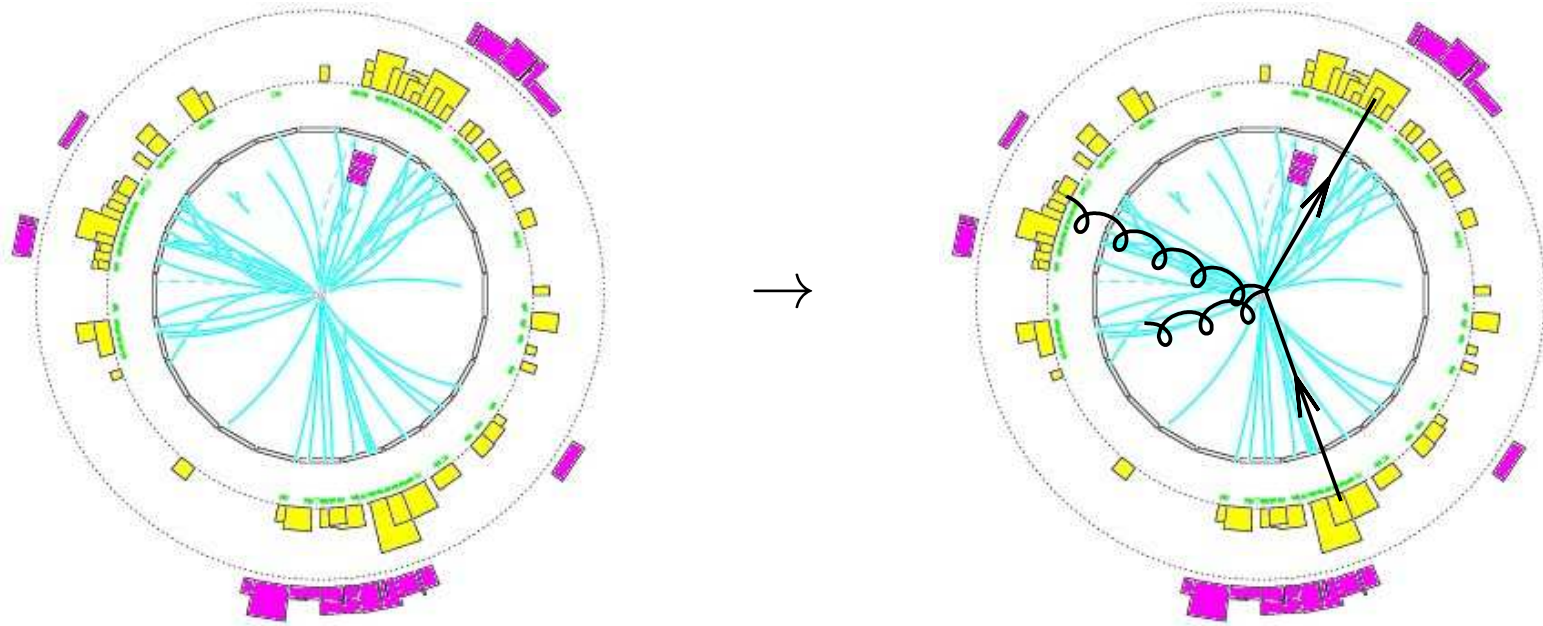


- “collinear” is arbitrary

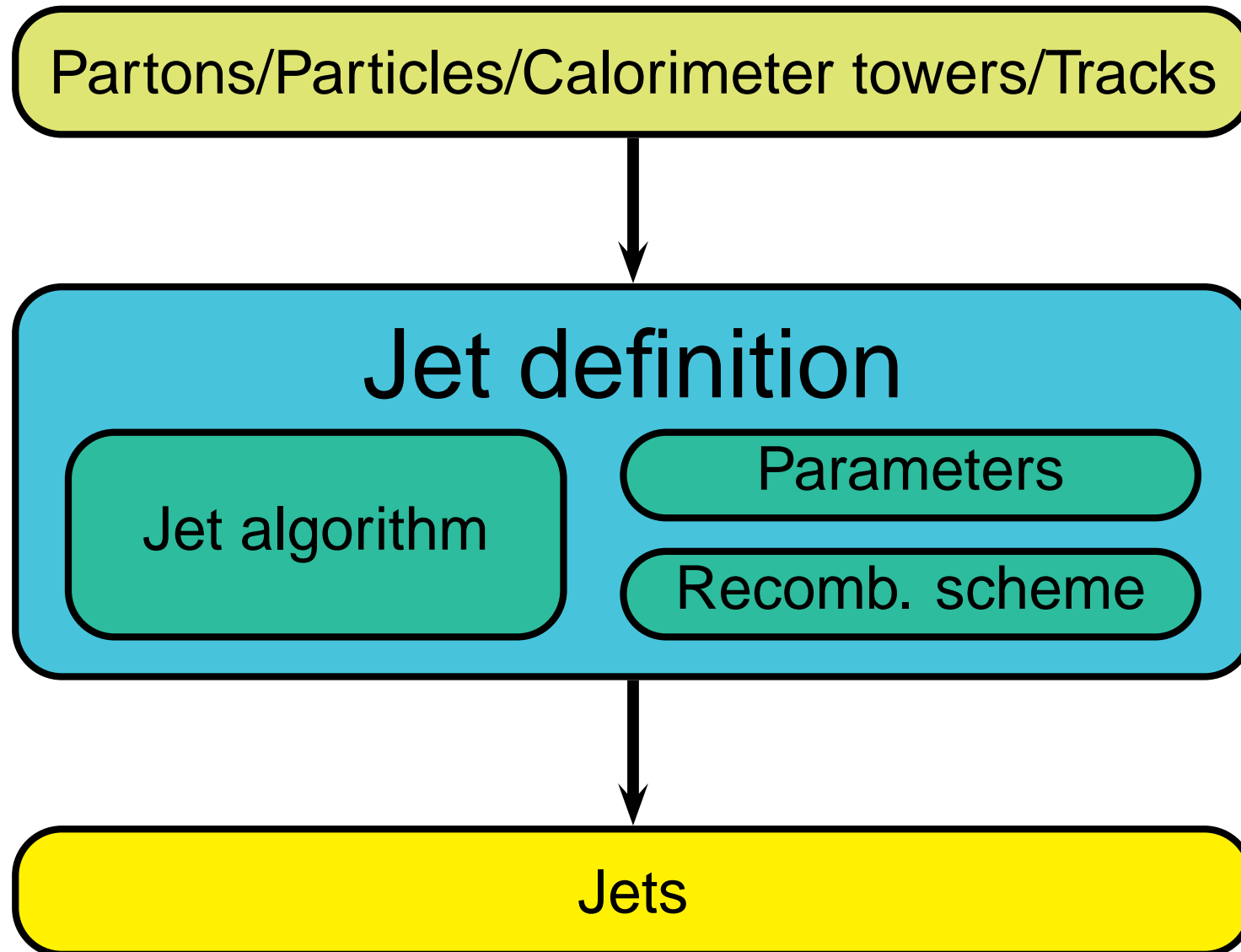
Jets and partons

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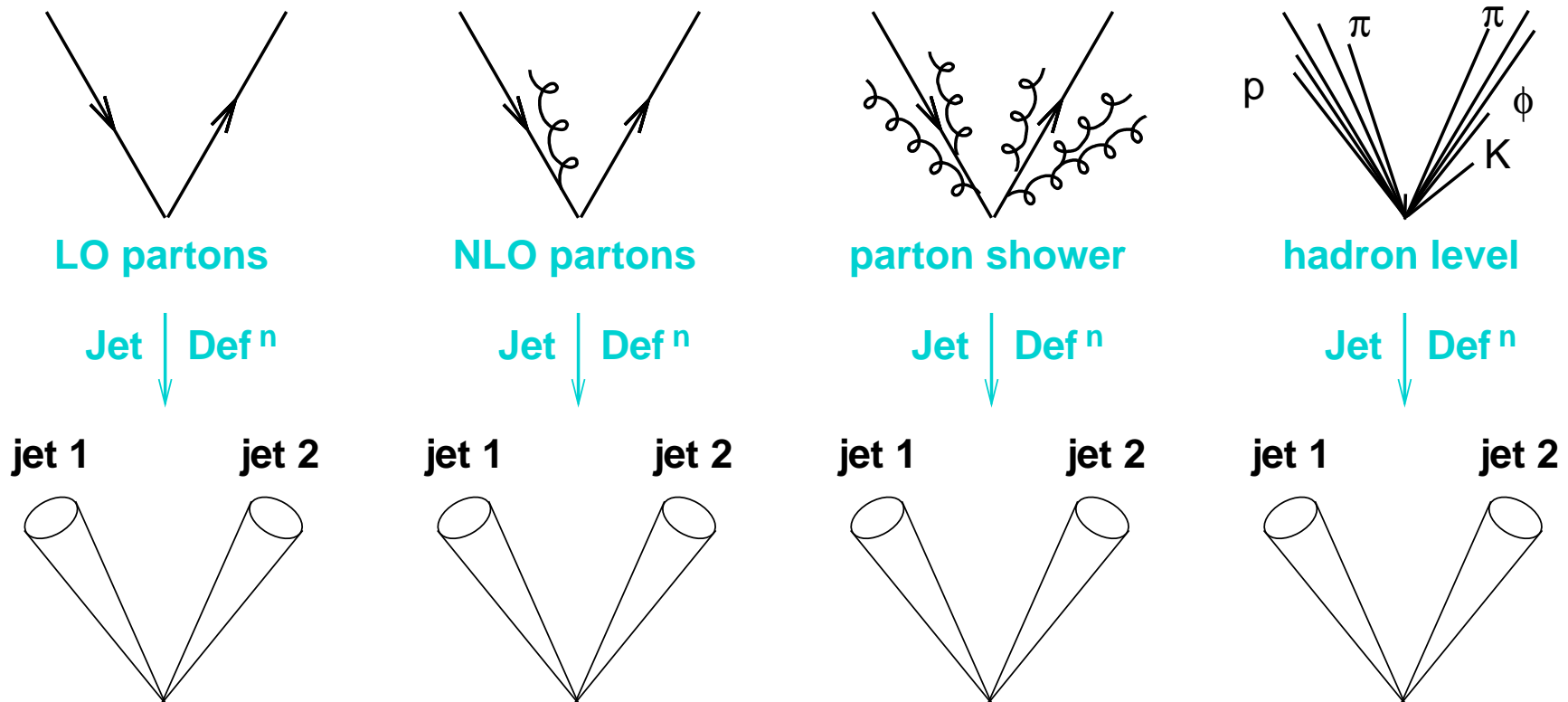
3 jets... or 4?



- “collinear” is arbitrary
- “parton” concept strictly valid only at LO



A jet definition is supposed to be (as) consistent (as possible) across different view of an event



Jet definitions: constraints

SNOWMASS accords (FermiLab, 1990)

Several important properties that should be met by a jet definition are [3]:

1. Simple to implement in an experimental analysis;
2. Simple to implement in the theoretical calculation;
3. Defined at any order of perturbation theory;
4. Yields finite cross section at any order of perturbation theory;
5. Yields a cross section that is relatively insensitive to hadronization.

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30 years later, these are only recently satisfied!!!

Proposal:
hire (many) PhD students
to look at the (many) millions of events

Proposal:

hire (many) ~~PhD students~~ grad students
to look at the (many) millions of events

Proposal:
hire (many) PhD students
to look at the (many) millions of events

Unless you have a better proposal?

Jet definition: successive recombinations

Idea: Undo the QCD cascade

- Define an inter-particle distance d_{ij} and a beam distance d_{iB}
- Successively
 - Find the minimum of all d_{ij}, d_{iB}
 - If d_{ij} , recombine $i + j \rightarrow k$ (remove i, j ; add k)
 - If d_{iB} , call i a jet (remove i)
- Until all particles have been clustered

Jet definition: successive recombinations

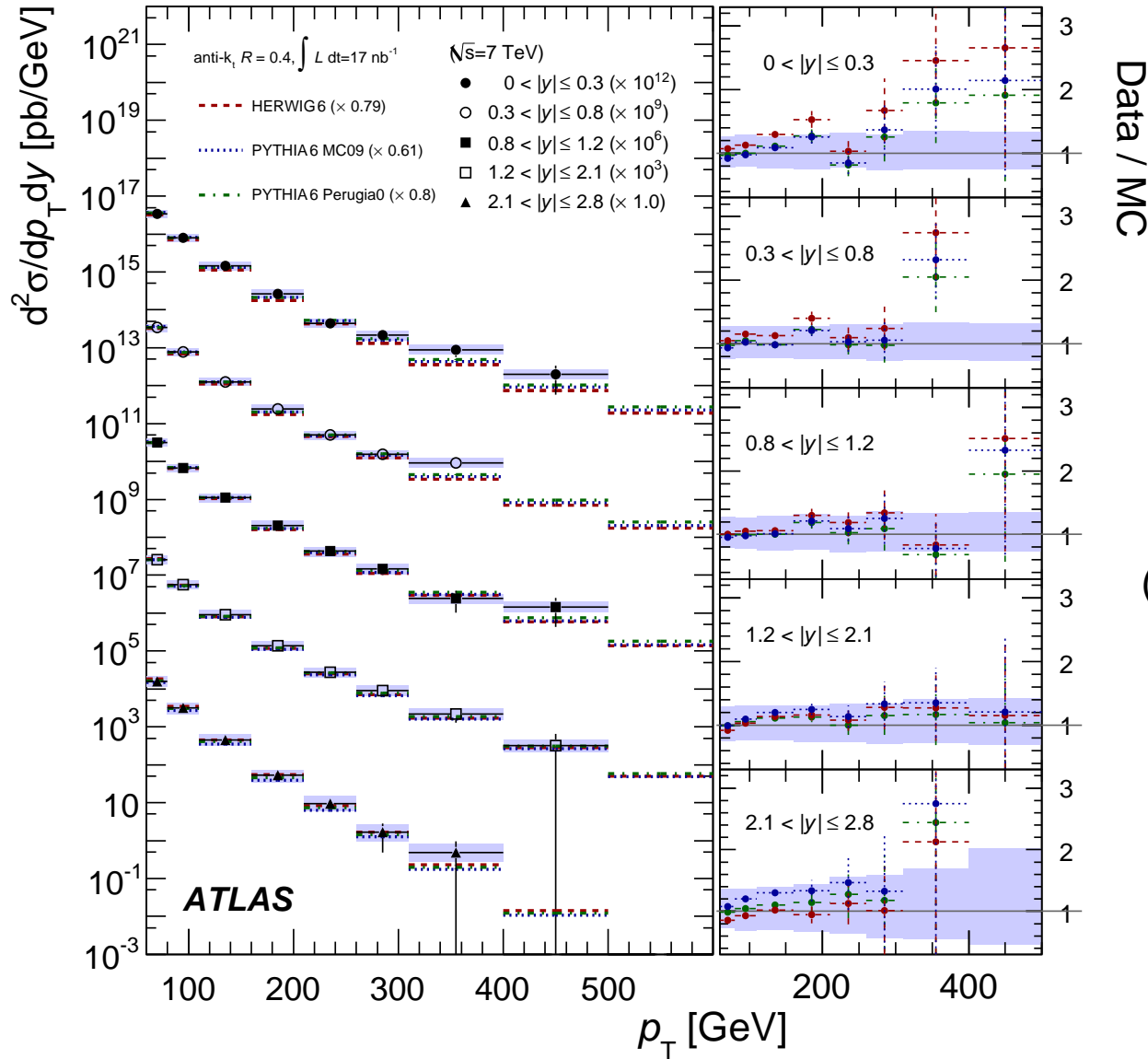
Typical choice of distances:

$$d_{ij}^2 = \min(k_{t,i}^{2p}, k_{t,j}^{2p})(\Delta y_{ij}^2 + \Delta\phi_{ij}^2)$$

$$d_{iB}^2 = k_{t,i}^{2p} R^2$$

- $p = 1$: k_t algorithm (1993)
(as close as possible to pQCD)
- $p = 0$: Cambridge-Aachen algorithm (1997)
(close to pQCD; useful for substructure)
- $p = -1$: anti- k_t algorithm (2008)
(circular/soft-resilient jets)

Jet definition at the LHC



Both CMS
and ATLAS
use the anti- k_t

(ATLAS: $R = 0.4$ and 0.6
CMS: $R = 0.5$ and 0.7)

Jets are alive

Still room for improvement:

- Experimentally:
jet energy scale
- Theoretically/Experimentally:
handle UE/pileup contamination
- Theoretically/Experimentally:
Tag boosted objects

Don't leave now...

...especially if you're on this list

- Sequential calibration (GSC) in ATLAS at the LHC
Reina CAMACHO
- Vers une mesure de la section efficace de production de paires des quarks top dans les canaux multileptons dans l'expérience ATLAS
Timothée THEVENEUX-PELZER
- Mesure de l'efficacité de l'étiquetage de jets beaux dans l'expérience ATLAS
Nancy TANNOURY
- Recherche de nouvelle physique avec ATLAS au LHC grâce à l'identification des jets de saveur b
Nicolas BOUSSON
- Four top events at the LHC from top-philic new physics
Léa GAUTHIER